# Integer Sequences from Queueing Theory 

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#### Abstract

Operators on probability distributions can be expressed as operators on the associated moment sequences, and so correspond to operators on integer sequences. Thus, there is an opportunity to apply each theory to the other. Moreover, probability models can be sources of integer sequences, both classical and new, as we show by considering the classical $M / G / 1$ single-server queueing model. We identify moment sequences that are integer sequences. We establish connections between the $M / M / 1$ busy period distribution and the Catalan and Schroeder numbers.


## 1 Introduction

Given the well established link between integer sequences and combinatorics, it is surprising that there are so few connections between integer sequences and queueing theory, because two prominent scholars, John Riordan (1902-1988) and Lajos Takács (1924- ), were active in both combinatorics and queueing theory [22, 23, 24, 28, 29, 30]. Of course, these authors saw connections; e.g., the index of Riordan's queueing book [23] contains entries for Bell polynomials, binomial moments, Catalan numbers, cumulant generating function, Lagrange expansion, and rooted trees.

In the OEIS [26], a search on "Riordan" yields 1221 entries, while a search on "queueing" produces 2 . We think Riordan would want less disparity. One exception is the entry submitted by A. Harel (A122525) making connection to the Erlang delay formula associated with the classical $M / M / s$ queue. A search on "queue" produces 15 entries, but these primarily
focus on combinatorial problems $[11,27]$ or a queue as a tool in computer science rather than queueing theory [12]; queueing theory concerns stochastic process models describing congestion, e.g., the probability distribution of customer waiting times [23, 28].

The purpose of this paper is to point out connections between the theories of probability and integer sequences, and to exhibit some integer sequences that arise in queueing theory. As a branch of probability theory, queueing theory has exploited the classical analytical approach to probability theory using transforms [23, 28], whose series expansions involve the integer moments of the probability distributions and close relatives. The entire probability distribution is characterized by the sequence of moments in great generality [10, 17]. Thus we were motivated to introduce an operational calculus for manipulating probability distributions on the positive halfline by either manipulating the associated Laplace transforms or the associated moment sequences [7]; a quick overview is provided by [7, Tables 1-3]. A more recent paper in the same spirit is [18].

The operational calculus for probability distributions on the positive halfline in the framework of moment sequences is closely related to operators commonly used to analyze integer sequences. Since many operators on integer sequences can be applied to moment sequences arising in probability theory, both when they are integers and when they are not, there is an opportunity for experts on integer sequences to contribute to probability theory through moment sequences. (The connection is also discussed in [13].) The probability connection also provides concrete models where integer sequences arise.

Here is how the present paper is organized. We start in $\S 2$ by reviewing moment sequences in probability theory and relating operators on probability distributions to operators on sequences. Then in $\S 3$ we review the classical $M / G / 1$ single-server queueing model and identify integer sequences arising in that context. We consider the $M / M / 1$ busy period distribution, its stationary excess distribution and the equilibrium time to emptiness. In that context we identify random quantities whose probability density functions have the Catalan and Large Schroeder numbers as moments. Finally, we consider the moment sequence for the $M / G / 1$ steady-state waiting time. We give the proofs of Theorems 5 and 11 in $\S 4$.

## 2 Moment Sequences in Probability

Let $Z$ be a nonnegative random variable with cumulative distribution function (cdf) $F$ and probability density function (pdf) f, i.e.,

$$
F(t) \equiv P(Z \leq t)=\int_{0}^{t} f(u) d u, \quad t \geq 0
$$

where $\equiv$ means "defined as." Let $\hat{f}(s)$ be the Laplace transform (LT) of $f$ (and thus $Z$ ) and let $\phi(x)$ be the associated moment generating function (mgf) of $f$, i.e.,

$$
\begin{align*}
\hat{f}(s) & \equiv \int_{0}^{\infty} e^{-s t} f(t) d t \equiv E\left[e^{-s Z}\right] \quad \text { and } \\
\phi(x) & \equiv \hat{f}(-x)=\int_{0}^{\infty} e^{x t} f(t) d t \equiv E\left[e^{x Z}\right] \tag{1}
\end{align*}
$$

where $s$ in (1) is understood to be a complex number with positive real part, while $x$ in (1) is a positive real number. The LT is always well defined; we assume that there exists $x^{*}>0$ such that $\phi(x)<\infty$ for $x<x^{*}$.

We obtain sequences by considering series expansions for the mgf $\phi$; in particular, we can write

$$
\begin{equation*}
\phi(x)=\sum_{n=0}^{\infty} \frac{m_{n} x^{n}}{n!}=\sum_{n=0}^{\infty} \mu_{n} x^{n} \tag{2}
\end{equation*}
$$

where of course we must have $\mu_{n}=m_{n} / n!, n \geq 0$. From the probability perspective, the object of primary interest is the cdf $F$, but the moments $m_{n}$ provide a useful partial characterization. The moment sequence can be calculated and employed to derive other quantities of interest in probability models $[1,7,16]$. There is a developing operational calculus for manipulating probability distributions via their moment sequences [7].

From the sequence perspective, we may instead regard the sequences $\left\{m_{n}: n \geq 1\right\}$ and $\left\{\mu_{n}: n \geq 1\right\}$ as the objects of primary interest. The two perspectives meet with the mgf $\phi$. From the sequence perspective, $\phi(x)$ arises as the generating function (gf) of the sequence $\left\{\mu_{n}: n \geq 1\right\}$, and we speak of $\left\{\mu_{n}\right\}$ as being its coefficients, while $\phi(x)$ arises as the exponential generating function (egf) of the sequence $\left\{m_{n}: n \geq 1\right\}$. Of course, $\left\{m_{n}: n \geq 1\right\}$ is always an integer sequence whenever $\left\{\mu_{n}: n \geq 1\right\}$ is, but not necessarily conversely. We will be giving examples where both are integer sequences.

A simple example from probability theory is the exponential distribution with mean $m$, specified by

$$
F(t) \equiv 1-e^{-t / m}, \quad f(t) \equiv(1 / m) e^{-t / m}, \quad t \geq 0, \quad \text { and } \quad \phi(x) \equiv(1-m x)^{-1}
$$

which has associated sequences

$$
m_{n} \equiv n!m^{n} \quad \text { and } \quad \mu_{n} \equiv m^{n}, \quad n \geq 0
$$

It is immediate that $\left\{\mu_{n}\right\}$ and $\left\{m_{n}\right\}$ are both elementary integer sequences whenever $m$ is an integer. Since the mean (first moment) is probabilistically only a scale parameter, depending on how the measurement units are defined, it is natural to follow the convention that the first moment is 1 ; here that gives $m_{1}=1$. Then we obtain the fundamental integer sequences $m_{n}=n!$ and $\mu_{n}=1, n \geq 1$.

As a consequence of the relations outlined above, we see that results about probability distributions can be translated into results about integer sequences, provided that the sequence $\left\{\mu_{n}\right\}$ or $\left\{m_{n}\right\}$ is indeed an integer sequence. Similarly, results about integer sequences can be translated into results about probability distributions, provided that the egf or gf of the integer sequence is indeed the mgf of a bonafide pdf.

We first observe that there is a natural probabilistic setting in which, not only is $\left\{m_{n}\right\}$ a moment sequence, but so is the associated sequence $\left\{\mu_{n}\right\}$. That occurs in spectral representations, where one pdf serves as a mixing pdf for another. In particular, suppose that the pdf $f$ can be represented as a continuous mixture of exponential pdf's via

$$
\begin{equation*}
f(t)=\int_{\tau_{1}}^{\tau_{2}} y^{-1} e^{-t / y} g(y) d y, \quad t \geq 0 \tag{3}
\end{equation*}
$$

in which case we call $g$ the mixing pdf for $f$ [4]. Append a superscript $f$ to $\left\{m_{n}\right\}$ and $\left\{\mu_{n}\right\}$ to denote dependence upon $f$. We observe that the sequence $\left\{\mu_{n}\right\}$ is itself the moment sequence of the associated mixing pdf $g$. Hence we call $\left\{\mu_{n}\right\}$ the mixing moments of $f$. (We omit the elementary proofs in this section.)

Proposition 1. (mixing moments) For pdf's $f$ and $g$ related via (3), $m_{n}^{g}=\mu_{n}^{f}, n \geq 1$.
A canonical operation in probability theory is convolution. If two independent nonnegative random variables $Z_{1}$ and $Z_{2}$ with pdf's $f_{Z_{1}}$ and $f_{Z_{2}}$ are added, then the sum $Z_{1}+Z_{2}$ has a pdf that is the convolution of the two component pdf's, i.e.,

$$
f_{Z_{1}+Z_{2}}(t)=\int_{0}^{t} f_{Z_{1}}(y) f_{Z_{2}}(t-y) d y, \quad t \geq 0
$$

One reason transforms are so frequently used in probability is that they convert convolution into a simple product; i.e., the associated mgf's are related by

$$
\phi_{Z_{1}+Z_{2}}(x)=\phi_{Z_{1}}(x) \phi_{Z_{2}}(x) .
$$

Because of the assumed stochastic independence, the moments are related by the binomial theorem,

$$
m_{n}^{Z_{1}+Z_{2}}=\sum_{k=1}^{n}\binom{n}{k} m_{k}^{Z_{1}} m_{n-k}^{Z_{2}} .
$$

Thus, if $\left\{m_{n}^{Z_{1}}\right\}$ and $\left\{m_{n}^{Z_{2}}\right\}$ are integer sequences, then so is $\left\{m_{n}^{Z_{1}+Z_{2}}\right\}$. For example, by above, that occurs when $Z_{1}$ and $Z_{2}$ have exponential or deterministic distributions with integer means.

With this background, we can interpret [7, Tables 1-3], which show various operators mapping one probability distribution into another. There are four columns: The first column contains the name and notation for the operator; the second column shows how the operator acts on LT's; the third column shows how it acts on pdf's; and the fourth column shows how it acts on moment sequences. From the perspective of integer sequences, the fourth column shows how it acts on the coefficients of the egf; we could then add a fifth column showing how it acts on the corresponding coefficients of the gf. Those familiar with integer sequences might want to translate the LT into the mgf by replacing $s$ with $-x$, and then interpret that mgf as the egf of the given $k^{\text {th }}$ moment.

In the rest of this section we highlight a few striking connections between operators on probability distributions, as in [7], and operators on integer sequences. First, a standard operator on integer sequences is the simple shift to the left, e.g., converting $1,1,2,6,22,90, \ldots$ to $1,2,6,22,90, \ldots$. We now show that, probabilistically, the simple shift applied to the coefficients $\mu_{n}$ corresponds to constructing the stationary-excess cdf of a cdf on the positive halfline having mean 1.

Given a nonnegative real-valued random variable $Z$ with cdf $F$ having finite moments $m_{k}, k \geq 1$, let $Z_{e}$ be a random variable with the associated stationary-excess cdf $F_{e}$ (a.k.a. the equilibrium excess or stationary residual-life cdf), defined by

$$
\begin{equation*}
F_{e}(t) \equiv P\left(Z_{e} \leq t\right) \equiv \frac{1}{m_{1}} \int_{0}^{t}(1-F(u)) d u, \quad t \geq 0 \tag{4}
\end{equation*}
$$

The stationary-excess cdf frequently arises in renewal theory; see of [25, Examples 7.16,7.17, $7.23,7.24]$ and [31]; it appears in [7, Table 2]. (A search on "renewal theory" in the OEIS gives two unrelated entries.) For us, the important fact is that the random variable $Z_{e}$ has moments

$$
\begin{equation*}
m_{e, k} \equiv E\left[Z_{e}^{k}\right]=\frac{m_{k+1}}{(k+1) m_{1}}, \quad k \geq 1 \tag{5}
\end{equation*}
$$

Hence, the transformation from a cdf of a nonnegative random variable to its associated stationary-excess cdf produces a simple shift on the gf coefficients.

Proposition 2. (the stationary-excess operator) Let $F$ be the $c d f$ of a nonnegative random variable with mean 1, mgf $\phi$ in (1) and associated sequence of mixing moments $\left\{\mu_{n}\right\}$ in (2) (coefficients of $\phi(x)$ when it is regarded as a $g f$ ). The associated stationary-excess cdf $F_{e}$ in (4) has mgf $\phi_{e}(x)=(\phi(x)-1) / x$. The mixing moments of $\phi_{e}(x)$ (coefficients of $\phi_{e}$ when it is regarded as a gf) are $\mu_{e, k}=\mu_{k+1}, k \geq 1$.

A similar relationship holds for the stationary-lifetime operator, mapping a pdf $f$ into the associated pdf

$$
\begin{equation*}
f_{s}(t) \equiv \frac{t f(t)}{m_{1}}, \quad t \geq 0 \tag{6}
\end{equation*}
$$

The stationary-lifetime pdf also frequently arises in renewal theory, e.g., [25, §7.7], and also appears in [7, Table 2]. For us, the important fact is that the moments are related by

$$
m_{s, k} \equiv \frac{m_{k+1}}{m_{1}}, \quad k \geq 1
$$

just like (5) without the $k+1$ in the denominator. Hence, the transformation from a pdf of a nonnegative random variable to its associated stationary-lifetime pdf produces a simple shift on the moments $m_{n}$ (the egf coefficients).

Proposition 3. (the stationary-lifetime operator) Let $f$ be the pdf of a nonnegative random variable with mean 1 , mgf $\phi$ in (1) and associated sequence of moments $\left\{m_{n}\right\}$ in (2). The associated stationary-lifetime pdf $f_{s}$ in (6) has mgf $\phi_{s}(x)=\phi^{\prime}(x)$. The coefficients of $\phi_{s}$ when it is regarded as an egf are $m_{s, k}=m_{k+1}, k \geq 1$.

We now turn to another basic probability operator, which is conveniently related to a continued fraction representation of the mgf $[8,9,13,15]$. If $\hat{f}$ is the LT of a pdf $f$, then the associated exponential mixture pdf has LT

$$
\begin{equation*}
\hat{f}_{\mathcal{E M}}(s) \equiv(1+s \hat{f}(s))^{-1} \tag{7}
\end{equation*}
$$

it appears in [7, Table 3]. The special case of exponential mixtures of inverse Gaussian (EMIG) distributions is discussed in $[7, \S 8]$ and in [9].

The probabilistic exponential mixing operator has a simple manifestation in the continued fraction representation of the mgf, regarding that mgf as a gf. Starting with the LT $\hat{f}(s)$ of a pdf $f$, if we represent the associated mgf as a formal power series by

$$
\begin{equation*}
\phi(x) \equiv \hat{f}(-x)=1+\mu_{1} x+\mu_{2} x^{2}+\mu_{3} x^{3}+\mu_{4} x^{4}+\mu_{5} x^{5}+\ldots, \tag{8}
\end{equation*}
$$

then the corresponding continued fraction ( CF ) is

$$
\begin{equation*}
\phi(x)=\frac{1}{1-} \frac{h_{1} x}{1-} \frac{h_{2} x}{1-} \frac{h_{3} x}{1-} \frac{h_{4} x}{1-} \ldots \tag{9}
\end{equation*}
$$

where $h_{1} \equiv \mu_{1}, h_{2} \equiv\left(\mu_{2}-\mu_{1}^{2}\right) / \mu_{1}$, etc. When $h_{n}>0$ for all $n$ we have an $S$-fraction and the underlying pdf $f$ is completely monotone (CM) [8].

Proposition 4. (the exponential-mixture operator) Let $f$ be the pdf of a nonnegative random variable with mean 1, mgf $\phi$ in (1) with associated sequence of CF coefficients $\left\{h_{n}\right\}$ in (9). The associated exponential-mixture pdf $f_{\mathcal{E} \mathcal{M}}$ with $L T$ in (7) has CF coefficients

$$
h_{\mathcal{E M}, 1}=1, \quad h_{\mathcal{E M}, n}=h_{n-1}, \quad n \geq 2
$$

i.e., it produces a shift to the right. The inverse exponential mixture operator [7, (7.3), p. 94] gives the corresponding shift to the left.

Proof. From the transform expression in (7), the conclusion is immediate: Given the CF representation for $\hat{f}(s)$, it is immediate that the corresponding CF for $(1+s \hat{f}(s))^{-1}$ shifts the coefficients one to the right; i.e., we write the CF for $(1+s \hat{f}(s))^{-1}$ as $1 /(1+s \hat{f}(s))$, inserting the CF for $\hat{f}(s)$.

## 3 Queueing Examples

### 3.1 The $M / G / 1$ Model

The $M / G / 1$ queue is a basic model in queueing theory, usually discussed in queueing textbooks; e.g., [25, §8.5]. There is a single server with unlimited waiting room. Customers arrive according to a Poisson process (the first $M$, for Markov) with rate $\lambda, 0<\lambda<\infty$. If the system is empty, then the customer goes immediately into service; otherwise the customer waits in queue. The successive service times come from a sequence of independent and identically distributed (i.i.d.) random variables with cdf $G$ having mean $1 / \mu, 0<\mu<\infty$. We will mostly consider the easiest case, in which the service-time cdf $G$ is exponential; then the model is denoted by $M / M / 1$.

A waiting customer enters service immediately upon service completion. Let $Q(t)$ be the number of customers in the system at time $t$ for $t \geq 0$. In the $M / M / 1$ model, the stochastic process $Q \equiv\{Q(t): t \geq 0\}$ is a birth-and-death stochastic process, with constant birth rate $\lambda$ and constant death rate $\mu$. Let $\rho \equiv \lambda / \mu$ be the traffic intensity. If $\rho<1$, then $P(Q(t)=j \mid Q(0)=i) \rightarrow(1-\rho) \rho^{j}$ as $t \rightarrow \infty$ for each $i$ and $j$; i.e., $Q(t)$ converges in distribution to a geometric distribution on the nonnegative integers, having mean $\rho /(1-\rho)$. We assume that $\rho<1$, under which the system is said to be stable. (If $\rho \geq 1$, then $P(Q(t) \leq j \mid Q(0)=i) \rightarrow 0$ as $t \rightarrow \infty$ for each $i$ and $j$.) For the $M / G / 1$ model, the steadystate distribution of $Q(t)$ is characterized by the Pollaczek-Khintchine transform [23, 28].

### 3.2 The $M / M / 1$ Busy Period

A busy period is the time from the arrival of a customer finding an empty system until the system is empty again $[23, \S 4.8]$. The first passage time of $Q$ from any state $j>0$ to $j-1$ is distributed as a busy period. Without loss of generality, we can measure time in units of mean service times, so that we let $\mu \equiv 1$. Then the model has the single parameter $\rho(=\lambda)$. Let $X_{\rho}$ denote a random variable with the busy period distribution, as a function of the traffic intensity $\rho$. Let $B_{\rho}$ be the cdf of $X_{\rho}$, i.e., $B_{\rho}(t) \equiv P\left(X_{\rho} \leq t\right), t \geq 0$, and let $b_{\rho}$ be the associated pdf. It turns out that

$$
b_{\rho}(t) \equiv \frac{1}{t \sqrt{\rho}} e^{-(1+\rho) t} I_{1}(2 t \sqrt{\rho}), \quad t \geq 0
$$

where $I_{1}(t)$ is the Bessel function of the first kind [23, (39), p. 63]. The LT of $b_{\rho}$ (and thus $X_{\rho}$ ) is

$$
\hat{b}_{\rho}(s) \equiv \int_{0}^{\infty} e^{-s t} b_{\rho}(t) d t \equiv E\left[e^{-s X_{\rho}}\right]=\frac{1+\rho+s-\sqrt{(1+\rho+s)^{2}-4 \rho}}{2 \rho},
$$

[23, (38), p. 63]. As usual, the moments can be obtained by differentiating the Laplace transform. The first two moments are $E\left[X_{\rho}\right]=(1-\rho)^{-1}$ and $E\left[X_{\rho}^{2}\right]=2(1-\rho)^{-3}$.

### 3.3 The Busy-Period Moment Sequence After Scaling

Our goal here is to obtain interesting integer sequences from the sequence of successive integer moments of a busy period $X_{\rho}$. To obtain integer sequences, we first perform a change of variables, introducing $\sigma \equiv \rho /(1-\rho)$ (the mean steady state number in system) or, equivalently, $\rho \equiv \sigma /(1+\sigma)$. Notice that the first two moments become $1+\sigma$ and $2(1+\sigma)^{3}$. It turns out that the entire moment sequence becomes an integer sequence whenever $\sigma$ is an integer. Hence, from this first step, we obtain an entire sequence of integer sequences.

To seek simple integer sequences, we further scale these moment sequences all to have mean (first moment) 1. We remark that this final spatial scaling also plays a role in understanding the performance of the $M / M / 1$ queue as the traffic intensity $\rho$ increases toward its critical value 1 . With appropriate scaling of both time and space, the stochastic process $Q$ approaches reflected Brownian motion with negative drift, while the busy period distribution has interesting behavior, in which both small values and large values play a role $[3,6,32]$.

In terms of random variables, let

$$
\begin{equation*}
Y_{\sigma} \equiv \frac{X_{\sigma /(1+\sigma)}}{1+\sigma}=(1-\rho) X_{\rho}, \quad \text { where } \quad \rho \equiv \frac{\sigma}{1+\sigma} \quad \text { and } \quad \sigma=\frac{\rho}{1-\rho} . \tag{10}
\end{equation*}
$$

Let $b(x ; \sigma)$ be the mgf of $Y_{\sigma}$. From above,

$$
\begin{equation*}
b(x ; \sigma) \equiv E\left[e^{x Y_{\sigma}}\right]=\frac{1+2 \sigma-x-\Psi(x)}{2 \sigma}, \quad \text { where } \quad \Psi(x) \equiv \sqrt{1-2(1+2 \sigma) x+x^{2}} \tag{11}
\end{equation*}
$$

Let the associated moments of $Y_{\sigma}$ be $m_{n}(\sigma) \equiv E\left[Y_{\sigma}^{n}\right], n \geq 1$, where $m_{1}(\sigma)=1$ for all $\sigma>0$. These moments, divided by $n$ !, are the coefficients of the series expansion of $b(x ; \sigma)$

$$
\begin{equation*}
b(x ; \sigma)=\sum_{n=0}^{\infty} \frac{m_{n}(\sigma) x^{n}}{n!}=\sum_{n=0}^{\infty} b_{n}(\sigma) x^{n}, \quad \text { where } \quad b_{n}(\sigma) \equiv \frac{m_{n}(\sigma)}{n!} \tag{12}
\end{equation*}
$$

The coefficients $b_{n}(\sigma)$ (and thus also the moments $m_{n}(\sigma)$ ) are polynomials in $\sigma$ (and thus integers when $\sigma$ is an integer); the first few are: $b_{0}(\sigma)=1, b_{1}(\sigma)=1, b_{2}(\sigma)=1+\sigma$ and $b_{3}(\sigma)=1+3 \sigma+2 \sigma^{2}$.

More specifically, we now relate the coefficients $b_{n}(\sigma)$ in (12) to the Catalan numbers, denoted by $C_{k}$, (A000108) starting with $1,1,2,5,14,42,132,429,1430$; in particular,

$$
\begin{equation*}
C_{n} \equiv \frac{1}{n+1}\binom{2 n}{n} \quad \text { and } \quad C_{n+1}=\sum_{i=0}^{n} C_{i} C_{n-i}, \quad n \geq 1 \tag{13}
\end{equation*}
$$

The Catalan numbers can be characterized in terms of their generating function

$$
\begin{equation*}
c(y) \equiv \sum_{k=0}^{\infty} C_{k} y^{k} \equiv \frac{1-\sqrt{1-4 y}}{2 y} \tag{14}
\end{equation*}
$$

Theorem 5. (mixing moments of the busy period pdf) For $n \geq 1$,

$$
\begin{equation*}
b_{n+1}(\sigma)=\sum_{k=0}^{n}\binom{n+k}{n-k} C_{k} \sigma^{k} \tag{15}
\end{equation*}
$$

where $b_{n}(\sigma)$ is defined in (11) and (12), and $C_{k}$ is the $k^{\text {th }}$ Catalan number in (13). In addition, there is a convenient recurrence relation, with $b_{0}(\sigma) \equiv b_{1}(\sigma) \equiv 1$ and

$$
\begin{equation*}
(n+1) b_{n+1}(\sigma)=(2 n-1)(1+\sigma) b_{n}(\sigma)-(n-2) b_{n-1}(\sigma) \tag{16}
\end{equation*}
$$

We provide a proof in $\S 4$. For $\sigma=1,2,3$, the coefficient sequences are:

$$
\begin{aligned}
& \left\{b_{n}(1)\right\}=1, \\
& \left\{b_{n}(2)\right\}= \\
& 1,
\end{aligned} 1, \quad 1, \quad 3, \quad 15, \quad 93, \quad 22, \quad 90, \quad 394, \quad(A 155069)
$$

Sequence $\left\{b_{n}(1)\right\}$ (A155069) is a relatively recent addition to the OEIS [26], having the title "Expansion of $\left(3-x-\sqrt{1-6 x+x^{2}}\right) / 2$ in powers of $x$." We provide a model context. However, sequence $\left\{b_{n}(1)\right\}$ shifted one to the left, i.e., $1,2,6,22,90, \ldots$ is (A006318), corresponding to the famous "Large Schroeder numbers." The "little Schroeder numbers" in (A001003) are obtained by dividing A006318 by 2 , i.e., the sequence $1,1,3,11,45, \ldots$.

The moment sequences themselves, $m_{n}(\sigma) \equiv n!b_{n}(\sigma)$ are of course also integer sequences, but they are evidently not in the OEIS [26]. For example, the sequences with terms $n!b_{n}(1)$ and $n!b_{n}(1 / 2), n \geq 0$, yielding $1,1,4,36,528,10800$ and $1,1,3,18,171,2250$, respectively, are not found.

The busy period is the first passage time of $Q$ from 1 to 0 . Since the first passage time pdf from each state $k$ to 0 is the $k$-fold convolution of the busy period pdf [3, Theorem 3.1]), the moments and mixing moments of these pdf's also generate integer sequences for integer $\sigma$.

Since we find regularity by multiplying $\left\{b_{n}(\sigma)\right\}$ by $1 / 2=\sigma /(1+\sigma)$ when $\sigma=1$, we are motivated to consider the associated sequences $\left\{\sigma b_{n+1}(\sigma) /(1+\sigma): n \geq 1\right\}$ for integers $\sigma>1$.

Corollary 6. For $n \geq 1$,

$$
\frac{\sigma b_{n+1}(\sigma)}{1+\sigma}=\sum_{k=0}^{n}(-1)^{n+k}\binom{n+k}{n-k} C_{k}(1+\sigma)^{k}
$$

We have already observed that Corollary 6 gives the little Schroeder numbers for $\sigma=1$; it is also discussed in [24, Problem 15, p. 168].

### 3.4 The Busy-Period Stationary-Excess Distribution

Let $B_{e}$ denote the busy period stationary-excess cdf associated with the busy-period cdf $B_{\rho}$, defined as in (4) after applying the scaling in (10). Let $b_{e}$ be the associated stationary-excess pdf. We can apply Proposition 2 to characterize the mixing moments $b_{e, n}(\sigma)$ of $b_{e}$.

Corollary 7. (mixing moments of the busy-period stationary-excess pdf) The $M / M / 1$ busyperiod stationary excess pdf $b_{e}$ has mgf

$$
\begin{align*}
b_{e}(x, \sigma) & \equiv \int_{0}^{\infty} e^{x t} b_{e}(t) d t \equiv \sum_{n=0}^{\infty} b_{e, n}(\sigma) x^{n}=\frac{b(x ; \sigma)-1}{x} \\
& =\frac{2}{1-x+\sqrt{1-2(1+2 \sigma) x+x^{2}}} \\
& =\frac{c\left(\sigma x /(1-x)^{2}\right)}{1-x}=\frac{2}{1-x+\Psi(x)} \tag{17}
\end{align*}
$$

where $c(y)$ is the generating function of the Catalan numbers in (14) and $\Psi$ is defined in (11). For $n \geq 1$,

$$
\begin{equation*}
b_{e, n}(\sigma)=b_{n+1}(\sigma), n \geq 1, \tag{18}
\end{equation*}
$$

for which an expression is given in Theorem 5. The mean of $b_{e}$ is $\sigma$.

### 3.5 Large Schroeder and Catalan Numbers as Moments of PDF's

In this section we identify the pdf's whose moments are (i) the Large Schroeder numbers and (ii) the Catalan numbers. By Corollary 7 and our observation after Theorem 5, the sequence $\left\{b_{e, n}(1)\right\}$ coincides with the Large Schroeder numbers (A006318). From [3, Theorem 5.1 and Corollary 5.2.1] and [4, Theorem 4.1], we can obtain the spectral representation for the pdf $b_{e}$. (We need to account for the different scaling of time used there.) That identifies the desired pdf and, by Proposition 1, the mixing moments.

Corollary 8. (Large Schroeder numbers as moments) The large Schroeder numbers arise as the moments of the mixing pdf associated with the stationary-excess busy-period pdf $b_{e}$ when $\sigma=1$. The mixing pdf for $b_{e}$ as a function of $\sigma$ is

$$
f(y ; \sigma)=\frac{\sqrt{(\tau-y)\left(y-\tau^{-1}\right)}}{2 \sigma \pi y}, \quad \frac{1}{\tau}<y<\tau
$$

where $\tau \equiv 1+2 \sigma+2 \sqrt{\sigma(1+\sigma)}$.

We introduced the Catalan numbers before Theorem 5. The Catalan numbers are related to reflected Brownian motion with drift -1 and diffusion coefficient 1 , denoted by $\{R(t)$ : $t \geq 0\}$. It is the limit of the $M / M / 1$ queue length process as $\rho \uparrow 1$ with appropriate scaling of time and space $[2,32]$. Since $E[R(t) \mid R(0)=0]$ is nondecreasing, $H_{1}(t) \equiv E[R(t) \mid R(0)=$ $0] / E[R(\infty)], t \geq 0$, is a cdf. By [2, Corollary 1.5.2], [7, §7] and [8, (8.7)], the pdf $h_{1}$ of $H_{1}$ has LT $\hat{h}_{1}(s)$, which can be characterized as the unique fixed point of the exponential mixture operator, i.e.,

$$
\begin{equation*}
\hat{h}_{1}(s)=\frac{1}{1+s \hat{h}_{1}(s)} . \tag{19}
\end{equation*}
$$

Together with the spectral representation given in [4, Theorem 4.1], that implies the following result. Without the probability model context, the result already appears in the commentary on (A000108) [21].
Theorem 9. (Catalan numbers as moments) The generating function $c(x)$ of the Catalan numbers in (13) and (14) coincides with the mgf of $h_{1}, \hat{h}_{1}(-x)$, associated with the LT in (19). As a consequence, the Catalan numbers arise as the moments of the mixing density of $h_{1}$,

$$
\begin{equation*}
f(y)=\frac{\sqrt{4-y}}{2 \pi \sqrt{y}}, \quad 0<y<4 \tag{20}
\end{equation*}
$$

If we take the Laplace transform of $f$ in (20) and replace $s$ by $-x$, then we obtain the mgf, which is the egf $\tilde{c}(x)$ of the Catalan numbers.
Corollary 10. The egf of the Catalan numbers is

$$
\tilde{c}(x) \equiv \sum_{n=0}^{\infty} \frac{C_{n} x^{n}}{n!}=e^{2 x}\left(I_{0}(2 x)-I_{1}(2 x)\right)
$$

where $I_{0}$ and $I_{1}$ are Bessel functions.

### 3.6 The $M / M / 1$ Equilibrium Time to Emptiness

From a sequence perspective, the shifting by 1 and multiplying by $\sigma /(1+\sigma)$ following Theorem 5 and Corollary 6 lead us to consider the mgf

$$
\begin{equation*}
p(x ; \sigma) \equiv \frac{1}{1+\sigma}+\frac{\sigma}{1+\sigma}\left(\frac{b(x ; \sigma)-1}{x}\right) . \tag{21}
\end{equation*}
$$

The function $p(x ; \sigma)$ in (21) turns out to be the moment generating function of the equilibrium time to emptiness in the $M / M / 1$ queue, i.e., the first passage time to state 0 by the stochastic process $Q$, assuming that $Q$ starts at time 0 according to its (geometric) steadystate distribution. Hence, with probability $(1-\rho)=1 /(1+\sigma)$ the process starts at 0 , so the first passage time is 0 . As a consequence, there is an atom at 0 with mass $1 /(1+\sigma)$. Conditional on starting at a positive value, the equilibrium time to emptiness coincides with the stationary-excess pdf $b_{e}$ of the busy-period pdf $b$ [5, Theorem 3]. Additional characterizations appear in [3, Theorem 3.3]. We make a connection to the exponential mixture operator in (7). Paralleling (12), we let $p_{n}(\sigma)$ be the coefficient of $p(x ; \sigma)$ as a gf, i.e., writing $p(x ; \sigma) \equiv \sum_{n=0}^{\infty} p_{n}(\sigma) x^{n}$. We provide a proof in $\S 4$.

Theorem 11. (equilibrium time to emptiness) The $m g f p(x ; \sigma)$ in (21) can be expressed as

$$
\begin{equation*}
p(x ; \sigma)=\frac{c\left((1+\sigma) x /(1+x)^{2}\right)}{1+x}=\frac{2}{1+x+\Psi(x)}, \tag{22}
\end{equation*}
$$

where $c(y)$ is the generating function of the Catalan numbers in (14) and $\Psi$ is defined in (11). Hence,

$$
\begin{equation*}
p_{n}(\sigma)=\frac{\sigma b_{n+1}(\sigma)}{1+\sigma}=\frac{\sigma}{1+\sigma} \sum_{k=0}^{n}\binom{n+k}{n-k} C_{k} \sigma^{k}, \quad n \geq 0 \tag{23}
\end{equation*}
$$

and the following recursion can be applied with $p_{0}(\sigma) \equiv 1$ and $p_{1}(\sigma) \equiv \sigma$

$$
\begin{equation*}
(n+2) p_{n+1}(\sigma)=(2 n+1)(1+2 \sigma) p_{n}(\sigma)-(n-1) p_{n-1}(\sigma), \quad n \geq 1 \tag{24}
\end{equation*}
$$

In addition, $b(x ; \sigma)$ can be expressed as an exponential mixture of $p(x ; \sigma)$, i.e.,

$$
\begin{equation*}
b(x ; \sigma)=\frac{1}{1-x p(x ; \sigma)} \tag{25}
\end{equation*}
$$

From (23) and (24), the sequence $\left\{p_{n}(\sigma): n \geq 1\right\}$ is an integer sequence for each positive integer $\sigma$. For $\sigma=1,2,3$, the coefficient sequences are:

$$
\begin{aligned}
& \left\{p_{n}(1)\right\}=1, \\
& \left\{p_{n}(2)\right\}=
\end{aligned} \quad 1, \quad 2, \quad 10, \quad 11, \quad 45, \quad 197, \quad(A 001003)
$$

We conclude this section by remarking that in the theory of integer sequences there also is a convenient invert operator, which can be expressed for LT's via

$$
\begin{align*}
\hat{g}(s) & \equiv \mathcal{I}(\hat{f}(s)) \equiv \operatorname{Excess}\left(\exp \operatorname{mixture}(\hat{f}(s))=\hat{f}_{\mathcal{E} \mathcal{M}, e}(s)\right. \\
& =\operatorname{Excess}\left(\frac{1}{1+s \hat{f}(s)}\right)=\frac{1}{s}\left(1-\frac{1}{1+s \hat{f}(s)}\right) \\
& =\frac{\hat{f}(s)}{1+s \hat{f}(s)} . \tag{26}
\end{align*}
$$

The following result links the three mgf's $b(x ; \sigma), b_{e}(x ; \sigma)$ and $(p(x ; \sigma)$; see Corollary 5.1 of [3]. We also consider the excess of the excess, denoted by $b_{e . e}$ [31].

Corollary 12. The mgf's $b_{e}(x ; \sigma), p(x ; \sigma)$ and $b(x ; \sigma)$ are related by

$$
b_{e}(x ; \sigma)=\mathcal{I}(p(x ; \sigma))=p(x ; \sigma) b(x ; \sigma)
$$

for the invert operator in (26) and

$$
b_{e, e}(x ; \sigma)=p_{e}(x ; \sigma) .
$$

Directly from the final expression in (17), we get a heavy-traffic limit for the scaled mgf with our scaling in (10), with the limit being the gf of the Catalan numbers. Let $X_{e}(\sigma)$ be a random variable with $\operatorname{mgf} b_{e}(x ; \sigma)$; i.e., $b_{e}(x ; \sigma) \equiv E\left[e^{x X_{e}(\sigma)}\right]$. Let $\Rightarrow$ denote convergence in distribution [32].

Corollary 13. As $\sigma \rightarrow \infty$,

$$
b_{e}(x / \sigma ; \sigma) \rightarrow \frac{2}{1-\sqrt{1-4 x}}=\frac{1-\sqrt{1-4 x}}{2 x} \equiv c(x)
$$

for $c$ in (14), so that

$$
\frac{X_{e}(\sigma)}{\sigma} \Rightarrow X \quad \text { as } \quad \sigma \rightarrow \infty
$$

where $c(x)=E\left[e^{x X}\right]$.
A related heavy-traffic limit for the time-scaled busy-period stationary-excess pdf was obtained in [6, Theorem 1]. (They are consistent.)

### 3.7 Continued Fractions and Hankel Transforms

We now look at the three probability mgf's $b(x ; \sigma), b_{e}(x, ; \sigma) \equiv(b(x ; \sigma)-1) / x$ and $p(x ; \sigma)$ in (11), (17) and (21) from the perspective of continued fractions and Hankel transforms. From the previous subsections, we know that the associated three random variables are closely related. That is shown again through this new perspective. First, the three mgf's can be represented as $S$ (Stieltjes) fractions. (In [8] we found that $S$ fractions frequently arise in CF's associated with birth-and-death stochastic processes.) The CF coefficients are given in Table 1.

| mgf | $h_{1}$ | $h_{2}$ | $h_{3}$ | $h_{4}$ | $h_{5}$ | $h_{6}$ |
| ---: | :--- | :--- | :--- | :--- | :--- | :--- |
| $b(x ; \sigma)$ | 1 | $\sigma$ | $1+\sigma$ | $\sigma$ | $1+\sigma$ | $\sigma$ |
| $b_{e}(x ; \sigma)$ | $1+\sigma$ | $\sigma$ | $1+\sigma$ | $\sigma$ | $1+\sigma$ | $\sigma$ |
| $p(x ; \sigma)$ | $\sigma$ | $1+\sigma$ | $\sigma$ | $1+\sigma$ | $\sigma$ | $1+\sigma$ |

Table 1: The coefficients of the continued fractions representing the three $M / M / 1 \mathrm{mgf}$ 's. The pattern repeats in each row.

For the generating functions of our $M / M / 1$ queueing examples, the determination of the CF coefficients is simple because we can apply the algebraic identity

$$
\frac{\alpha}{1+} \frac{\gamma}{1+} \frac{\alpha}{1+} \frac{\gamma}{1+} \ldots=\frac{1}{2}\left(\sqrt{1+2(\gamma+\alpha)+(\gamma-\alpha)^{2}}-1-(\gamma-\alpha)\right)
$$

for constants $\alpha$ and $\gamma$; i.e., we can solve the equation

$$
\zeta=\frac{\alpha}{1+} \frac{\gamma}{1+\zeta}
$$

for $\zeta$.
The Hankel transform of an integer sequence provides a useful partial characterization; it is a many-to-one function mapping an integer sequence into another integer sequence. (For an example of non-uniqueness, see Corollary 14 below.) For example, the Hankel transform of the Catalan numbers is the sequence $1,1,1, \ldots$ [13, 20]. Following [15, 17], starting from a sequence $\left\{\omega_{n}: n \geq 0\right\} \equiv \omega_{0}, \omega_{1}, \omega_{2}, \omega_{2}, \ldots$ with $\omega_{0} \equiv 1$, let the Hankel matrix $M^{(n)}$ be the $(n+1) \times(n+1)$ symmetric matrix with elements $M_{i, j}^{(n)} \equiv \omega_{i+j-2}, 0 \leq i \leq n, 0 \leq j \leq n$. (The first row contains the first $n+1$ elements and $M_{n+1, n+1} \equiv \omega_{2 n}$. Let $H_{2 n} \equiv \operatorname{det}\left(M^{(n)}\right)$, the even Hankel determinant. Let the Hankel transform of the sequence $\left\{\omega_{n}: n \geq 0\right\}$ above be the sequence $\left\{H_{2 n}: n \geq 0\right\}$; it starts with $H_{0}=1$.

One way to compute the Hankel transform of a sequence $\left\{\omega_{n}: n \geq 0\right\}$ is to first determine the corresponding continued fraction from the formal power series; i.e., starting from (8), we determine (9) above. The coefficients $h_{n}$ appearing in (9) can be determined by applying the normalized Viskovatov algorithm, as given on [15, p. 112]. Then we apply the iteration

$$
H_{2 n}=\left(\prod_{i=1}^{2 n} h_{i}\right) H_{2 n-2}, \quad n \geq 1
$$

see Theorem 1.4.10 on of [17, p. 23]. Another way to arrive at this result is via the even contraction of the CF in (9), as given in [13, (12.3), p. 270]; note that $\beta_{n}$ in [13] is equal to $h_{2 n-1} h_{2 n}, n \geq 1$, in our notation. Then (12.2) of [13, (12.2)] and our iteration yield the same result.

Corollary 14. The Hankel transforms of the sequences $\left\{b_{n}(\sigma)\right\}$ in Theorem $5,\left\{b_{e, n}(\sigma)\right\}$ in Corollary 7 and $\left\{p_{n}(\sigma\}\right.$ in Theorem 11 are

$$
\begin{aligned}
& H_{2 n}(b)=\sigma^{\nu(n)}(1+\sigma)^{\nu(n)-n}, \\
& H_{2 n}(p)=H_{2 n}\left(b_{e}\right)=\left(\sigma+\sigma^{2}\right)^{\nu(n)},
\end{aligned}
$$

where $\nu(n) \equiv n(n+1) / 2$.

### 3.8 The Stationary Waiting Time in the $M / G / 1$ Queue

The mgf $w(x ; \sigma) \equiv E\left[e^{x W}\right]$ of the steady-state waiting time $W$ in the $M / G / 1$ queue can be expressed in terms of the mgf $g_{e}(x ; \sigma)$ of the stationary-excess of the general service-time distribution using the Pollazcek-Khintchine transform (again assuming that the mean service time is 1 ) by applying the random sum operator from [7, Table 1], yielding

$$
\begin{equation*}
w(x ; \sigma)=\frac{1-\rho}{1-\rho g_{e}(x ; \sigma)}=\frac{1}{1+\sigma-g_{e}(x ; \sigma)} \tag{27}
\end{equation*}
$$

[23, §4.3] and [5]. Riordan [23, (18a), p. 49] and Takács [28] develop a nice recursion for the moments, which as a function of $\sigma$ becomes

$$
\begin{equation*}
E\left[W^{n-1}(\sigma)\right]=\frac{\sigma}{n} \sum_{k=2}^{n}\binom{n}{2} g_{k} E\left[W^{n-k}(\sigma)\right] \quad \text { and } \quad E\left[W^{0}(\sigma)\right] \equiv 1 \tag{28}
\end{equation*}
$$

where $g_{k}$ is the $k^{\text {th }}$ moment of the service-time cdf. Clearly, (28) can be the source of many integer sequences.

To illustrate, we now consider the Catalan numbers as service-time moments, which is legitimate by Theorem 9. We exploit the following result, obtained by combining [2, Corollaries 1.3.2 and 1.5.1] with Theorem 9.

Lemma 15. Let $h_{1}$ be the density of the RBM first moment function, which has the Catalan numbers as its mixing moments. Then

$$
\begin{equation*}
\hat{h}_{1, e} s=\hat{h}_{1}(s)^{2}, \quad \text { so that } \quad c_{e}(x)=c(x)^{2} \tag{29}
\end{equation*}
$$

where $c(x)$ is the generating function of the Catalan numbers.
Theorem 16. If the service time pdf is the RBM first moment pdf $h_{1}$ characterized by its $L T$ in (19), which has the Catalan numbers as its moments, then the waiting time mgf in (27) becomes

$$
\begin{equation*}
w(x ; \sigma)=\frac{1}{1+\sigma-\sigma c(x)^{2}} \tag{30}
\end{equation*}
$$

and the moments of the stationary waiting time pdf are

$$
\begin{equation*}
E\left[W^{n-1}(\sigma)\right]=\sigma \sum_{k=2}^{n} \frac{(n-1)!}{(n-k)!} C_{k} E\left[W^{n-k}(\sigma)\right] \tag{31}
\end{equation*}
$$

Hence, if $\sigma$ is an integer, then $\left\{E\left[W^{n}(\sigma)\right]\right\}$ is an integer sequence.
For the case $\sigma=1$, we find $\left\{E\left[W^{n}(1)\right]\right\}=1,2,18,252,4776, \ldots$, which is not in the OEIS. Riordan [23, (19), p. 50] shows how the moments $E\left[W^{n}(\sigma)\right]$ can be expressed in terms of the multivariate Bell polynomials.

Let $w_{n}(\sigma) \equiv E\left[W^{n}(\sigma)\right] / n$ ! be the associated mixing moments. From (31), we obtain a recursion for $w_{n}(\sigma)$.

Corollary 17. A recursion for the mixing moments $w_{n}(\sigma)$ defined above is

$$
\begin{equation*}
w_{n-1}(\sigma)=\sigma \sum_{k=2}^{n} C_{k} w_{n-k}(\sigma), \quad n \geq 1 \tag{32}
\end{equation*}
$$

For $\sigma=1$, we get $\left\{w_{n}(1)\right\}=1,2,9,42,199, \ldots$, which also is not yet in OEIS.

### 3.9 General Birth-and-Death Processes

The most relevant previous work connecting queueing theory to integer sequences and providing a suitable framework for generalization evidently is the previous work connecting birth-and-death processes to continued fractions; see [8, 18] and references therein; continued fractions are known to be intimately connected to integer sequences. We have already mentioned that the stochastic process $Q \equiv\{Q(t): t \geq 0\}$, representing the number of customers in an $M / M / 1$ queueing model at time $t$, is a birth-and-death stochastic process.

Indeed, it is a special birth-and-death process with constant birth rate $\lambda$ and constant death rate $\mu$. More generally, these rates are functions of the state [25, $\S 6.3] ; \lambda_{k}$ is the rate up and $\mu_{k}$ is the rate down when $Q(t)=k$. The general birth-and-death process represents a more general queueing model.

For a general birth-and-death process $X \equiv\{X(t): t \geq 0\}$, an important quantity is the probability

$$
\begin{equation*}
P_{0,0}(t) \equiv P(X(t)=0 \mid X(0)=0), \quad t \geq 0 ; \tag{33}
\end{equation*}
$$

see $[3, \S 11],[18, \S 3.1]$ and [19]. From [18, (4.1), p. 771], the LT of $P_{0,0}(t), \hat{P}_{0,0}(s)$, has a representation as an $S$ fraction

$$
\begin{align*}
\hat{P}_{0,0}(s) & =\frac{1}{s+} \frac{\lambda_{0}}{1+} \frac{\mu_{1}}{s+} \frac{\lambda_{1}}{1+} \frac{\mu_{2}}{s+} \ldots \\
& =\frac{z}{1+} \frac{\lambda_{0} z}{1+} \frac{\mu_{1} z}{1+} \frac{\lambda_{1} z}{1+} \frac{\mu_{2} z}{1+} \ldots \tag{34}
\end{align*}
$$

where $z \equiv 1 / s$. (The last line follows from $[8,(1.5)]$ using the sequence $\left\{c_{n}: n \geq 0\right\}$ with $c_{2 n} \equiv 1$ and $c_{2 n+1} \equiv z, n \geq 1$.) We can find the corresponding power series

$$
\hat{P}_{0,0}(s)=z\left(1-p_{1} z+p_{2} z^{2}-p_{3} z^{3}+\ldots\right) ;
$$

see $\left[8,(3.9),(3.10)\right.$, p. 397]. The sequence $\left\{1, p_{1}, p_{2}, p_{3}, \ldots\right\}$ may be an integer sequence.
For the special case of the $M / M / 1$ queue, with our scaling we have $\lambda_{i}=\sigma$ and $\mu_{i+1}=1+\sigma$ for all $i, i \geq 0$, so that

$$
\hat{P}_{0,0}(s)=z\left(\frac{1}{1+} \frac{\sigma z}{1+} \frac{(1+\sigma) z}{1+} \frac{\sigma z}{1+} \frac{(1+\sigma) z}{1+} \ldots\right) .
$$

Then from Table 1 we have the relation

$$
\begin{equation*}
\hat{P}_{0,0}(s)=z \hat{p}(z) \tag{35}
\end{equation*}
$$

where $\hat{p}(s)$ is the Laplace transform of the equilibrium time to emptiness.

## 4 Proofs

In this concluding section we provide proofs for Theorems 5 and 11.

### 4.1 Proof of Theorem 5.

We give two proofs. The first is direct; the second applies [14].
Direct proof of (15). From [29, pp. 232-233], after a change of scale and notation, we have

$$
b_{n+1}=\left(\frac{1+\sigma}{n+1}\right) \sum_{k=1}^{n}\binom{n+k}{k}\binom{n-1}{k-1} \sigma^{k-1}
$$

(This equation is also given by $[24,(21)$, p. 151] after we make the identification that $H_{n}(z)=p_{n}(\sigma)=\sigma b_{n+1}(\sigma) /(1+\sigma)$.) On the right side, move $(1+\sigma)$ inside the sum and identify the coefficients of $\sigma^{k}$, obtaining

$$
\left(\frac{1}{n+1}\right)\left(\binom{n+k+1}{k+1}\binom{n-1}{k}+\binom{n+k}{k}\binom{n-1}{k-1}\right)=\binom{n+k}{n-k}\binom{2 k}{k}\left(\frac{1}{k+1}\right)
$$

Direct proof of (16). We follow the technique of [23, p. 107] to establish a three-term recursion. By the successive differentiation with respect to $x$ in (11), we find that

$$
\Psi(x)^{2} b^{\prime \prime}(x ; \sigma)=(1+\sigma)(1+2 \sigma-x) b^{\prime}(x ; \sigma)+(1+\sigma)^{2} b(x ; \sigma)
$$

In this equation, make the following substitutions:

$$
\begin{array}{r}
b(x ; \sigma)=\sum_{n=0}^{\infty} b_{n}(\sigma) x^{n}, \\
b^{\prime}(x ; \sigma)=\sum_{n=0}^{\infty}(n+1) b_{n+1}(\sigma) x^{n} \\
b^{\prime \prime}(x ; \sigma)=\sum_{n=0}^{\infty}(n+1)(n+2) b_{n+2}(\sigma) x^{n} .
\end{array}
$$

Then collect and equate the coefficients of $x^{n-1}$ and the result follows.
Application of [14]. The idea in this second proof is to directly make connection to the Catalan numbers via their generating function $c$ in (14). Starting from (11), we apply (17), (21) and the proof of (22) in Theorem 11 below to conclude that

$$
\begin{equation*}
b_{e}(x ; \sigma)=\frac{1}{1-x} c\left(\frac{\sigma x}{(1-x)^{2}}\right) \quad \text { and } \quad p(x ; \sigma)=\frac{1}{1+x} c\left(\frac{(1+\sigma) x}{(1+x)^{2}}\right) \tag{36}
\end{equation*}
$$

We then apply [14, Proposition 15, p.12] (with change of variables $y \equiv \sigma x$ and $y \equiv(1+\sigma) x$, respectively) to immediately deduce the conclusions

$$
\begin{equation*}
b_{e, n}(\sigma)=\sum_{k=0}^{n}\binom{n+k}{n-k} C_{k} \sigma^{k} \quad \text { and } \quad p_{n}(\sigma)=\sum_{k=0}^{n}\binom{n+k}{n-k}(-1)^{n+k} C_{k}(1+\sigma)^{k} \tag{37}
\end{equation*}
$$

given in Corollary 7 and Theorem 11. We then can apply (18) and (17) to get the conclusion in Theorem 5 for $b$ from $b_{e}$. By this line of reasoning, we establish Theorems 5 and 11 and Corollaries 6 and 7 all at once, and have a new perspective on their relationship.

### 4.2 Proof of Theorem 11.

Proof of (22). From the representation of $p(x ; \sigma)$ in (21), we have

$$
\begin{aligned}
p(x ; \sigma) & =\frac{1}{1+\sigma}+\frac{\sigma}{1+\sigma}\left(\frac{2}{1-x+\Psi(x)}\right) \\
& =\frac{1}{1+\sigma}+\frac{\sigma}{1+\sigma}\left(\frac{1-x-\Psi(x)}{2 \sigma x}\right) \\
& =\frac{1+x-\Psi(x)}{2(1+\sigma) x}=\frac{2}{1+x+\Psi(x)}
\end{aligned}
$$

On the other hand, by (14),

$$
\begin{aligned}
\frac{c\left((1+\sigma) x /(1+x)^{2}\right)}{1+x} & =\frac{2}{1+x+(1+x) \sqrt{1-4(1+\sigma) x /(1+x)^{2}}} \\
& =\frac{2}{1+x+\Psi(x)}
\end{aligned}
$$

Proof of (23). Combine (15) and (21). ■
Proof of (24). Combine (16) and (21). ■
Proof of (25). From (11),

$$
\begin{aligned}
b(x ; \sigma) & =\frac{1+2 \sigma-x-\Psi(x)}{2 \sigma}=\frac{2(1+\sigma)}{1+2 \sigma-x+\Psi(x)} \\
& =\frac{1}{1-x([1+x-\Psi(x)] /[2(1+\sigma) x])}=\frac{1}{1-x p(x ; \sigma)}
\end{aligned}
$$

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