

A Generalized Apéry Series

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Abstract

The inverse central binomial series

$$S_k(z) = \sum_{n=1}^{\infty} \frac{n^k z^n}{\binom{2n}{n}}$$

popularized by Apéry and Lehmer is evaluated for positive integers k along with the asymptotic behavior for large k . It is found that value $z = 2$, as commented on by D. H. Lehmer provides a unique relation to π .

1 Introduction

Since the appearance of $S_{-3}(1)$ in Apéry's famous proof [1] in 1979 that $\zeta(3)$ is irrational, an extensive literature has been devoted to the series

$$S_k(z) = \sum_{n=1}^{\infty} \frac{n^k z^n}{\binom{2n}{n}} \quad (1)$$

For example, in 1985 Lehmer [2] presented a number of special cases which could be obtained from the Taylor series for $f(x) = x^{-1/2}(1-x)^{-1/2}\sin^{-1}x$ using only elementary calculus. In passing, he noted that when k is a positive integer, $S_k(2)$ had the form $a_k - b_k\pi$ and that the rational number a_k/b_k "is a close approximation to π . This remark was recently taken up by Dyson et al. [3], who proved that $|a_k/b_k - \pi| = O(Q^{-k})$ as $k \rightarrow \infty$ where $Q = \sqrt{1 + (2\pi/\ln 2)^2}$. Lehmer also showed that for positive integer k

$$S_k(z) = \frac{2^{k+z^{5/2}} z^{1/2}}{(4-z)^{k+3/2}} (A_k(z/4) \sin^{-1}(\sqrt{z/4}) + \sqrt{z(4-z)} B_k(z/4)) \quad (2)$$

where A_k and B_k are recursively defined polynomials. It was apparently not until 2005 that (2) was evaluated explicitly, for $z = 1$, by J. Borwein and P. Girgensohn [4] who showed

$$S_k(1) = \frac{1}{2} (-1)^{k+1} \sum_{j=1}^{k+1} (-1)^j j! S(j+1, j) 3^{-j} \binom{2j}{j} \left(\sum_{i=1}^{j-1} \frac{3^i}{(2i+1) \binom{2i}{i}} + \frac{2}{3\sqrt{3}} \pi \right). \quad (3)$$

where the *Stirling numbers of the second kind* are defined by

$$S(k, j) = \frac{(-1)^j}{j!} \sum_{m=0}^j (-1)^m m^k \binom{j}{m}. \quad (4)$$

The aim of the present note is to extend (3) to complex z and thus to continue (1) analytically beyond its circle of convergence $|z| = 4$.

2 Calculation

We begin with the observation that $(m \binom{2m}{m})^{-1} = B(m, m+1)$, where B denotes Euler's beta integral. Hence,

$$S_k(z) = \int_0^1 \frac{dt}{t} \sum_{m=1}^{\infty} m^{k+1} (zt(1-t))^m. \quad (5)$$

Next, equation (21) of Girgensohn and Borwein [4],

$$\sum_{m=1}^{\infty} m^p X^m = \sum_{n=1}^p \sum_{m=1}^n (-1)^{m+n} \binom{n}{m} m^p X^n (1-X)^{-n-1}, \quad (6)$$

gives

$$S_k(z) = \sum_{n=1}^{k+1} \sum_{m=1}^n (-1)^{m+n} \binom{n}{m} m^{k+1} \int_0^1 \frac{dt}{t} \frac{(zt(1-t))^n}{(1-zt(1-t))^{n+1}}. \quad (7)$$

In the appendix it is shown that

$$\int_0^1 \frac{dt}{t} \frac{(zt(1-t))^n}{(1-zt(1-t))^{n+1}} = \frac{\sqrt{\pi}\Gamma(n)}{\Gamma(n+1/2)} X^n {}_2F_1(-1/2, n; n+1/2; -X) \quad (8)$$

where $X = z/(4-z)$, so

$$S_k(z) = \sum_{n=1}^{k+1} n! B(n, 1/2) S(k+1, n) X^n {}_2F_1(-1/2, n; n+1/2; -X). \quad (9)$$

By induction, starting with the tabulated value for $n = 1$ and using Gauss' contiguity relations we find (some details are given in the appendix)

$$\begin{aligned} & {}_2F_1(-1/2, n; n+1/2; -X) = \\ & \left(\frac{1}{2}\right)_n \left(\frac{1}{n!} + \frac{1}{\sqrt{\pi}\Gamma(n)} \sum_{k=0}^{n-1} \frac{(-1)^k \Gamma(k+1/2)}{(k+1)!} \binom{n-1}{k} \left(\frac{X+1}{X}\right)^{k+1} \times \right. \\ & \quad \left. \left[\sqrt{X} \sin^{-1} \sqrt{\frac{X}{X+1}} - \frac{1}{2} \sum_{l=1}^k \frac{(l-1)!}{(1/2)_l} \left(\frac{X}{X+1}\right)^l \right] \right). \end{aligned} \quad (10)$$

(We have used the ascending factorial notation $(a)_n = \Gamma(a+n)/\Gamma(a)$). Therefore we have the principal result

$$S_k(z) = \sum_{n=1}^{k+1} n! \left(\frac{z}{4-z} \right)^n S(k+1, n) \times \left(\frac{1}{n} + \sum_{p=0}^{n-1} (-1)^p \frac{(1/2)_p}{(p+1)!} \binom{n-1}{p} \left(\frac{4}{z} \right)^{p+1} \left(\sqrt{\frac{z}{4-z}} \sin^{-1} \frac{\sqrt{z}}{2} - \frac{1}{2} \sum_{l=1}^p \frac{\Gamma(l)}{(1/2)_l} \left(\frac{z}{4} \right)^l \right) \right) \quad (11)$$

Equation (11) is rather condensed; in unpacking it, sums with upper limit less than the lower limit are to be interpreted as 0. It is clear from (11) that for *rational* z

$$\sum_{m=1}^{\infty} \frac{n^k z^n}{\binom{2n}{n}} = R_1(z, k) + R_2(z, k) \sqrt{\frac{z}{4-z}} \sin^{-1} \frac{\sqrt{z}}{2}, \quad (12)$$

where R_j is a rational number.

One sees from (11) that $S_k(z)$ is analytic on the two-sheeted Riemann surface formed by two planes cut and rejoined along the real half-line $x > 4$. The numbers in (12) have the explicit expressions

$$R_1(z, k) = \quad (13)$$

$$\sum_{n=1}^{k+1} n! S(k+1, n) \left(\frac{z}{4-z} \right)^n \left(\frac{1}{n} - \frac{1}{2} \sum_{p=1}^{n-1} \sum_{l=1}^p \frac{(-1)^p (1/2)_p}{(p+1)! (1/2)_l} \binom{n-1}{p} \Gamma(l) \left(\frac{4}{z} \right)^{p-l+1} \right),$$

$$R_2(z, k) = \sum_{n=1}^{k+1} n! S(k+1, n) \sum_{p=0}^{n-1} \frac{(-1)^p}{(p+1)!} \binom{n-1}{p} \left(\frac{4}{z} \right)^{p+1}. \quad (14)$$

3 Asymptotics

It is convenient to work in terms of the exponential generating function

$$G(z, t) := \sum_{k=0}^{\infty} S_k(z) \frac{t^k}{k!} = S_0(z e^t) = \frac{z}{4 - z e^t} + \frac{4\sqrt{z} e^{t/2}}{(4 - z e^t)^{3/2}} \sin^{-1} \frac{\sqrt{z} e^{t/2}}{2} \quad (15)$$

To find the generating functions $\rho_j(z, t) := \sum R_j(z, k) t^k / k!$, it would be simplest to start with a series $D_k(z) = R_1(z, k) - R_2(z, k) \sqrt{\frac{z}{4-z}} \sin^{-1} \frac{\sqrt{z}}{2}$, work out its generating function $D(z, t)$ and by taking the sum and difference identify ρ_1 and ρ_2 . However, this series has not been found and there is nothing to guarantee its existence in tractable form. Therefore, the ρ_j were evaluated directly from (13) and (14). The details are omitted as the results

$$\rho_1(z, t) = \frac{z e^t}{4 - z e^t} + \frac{8}{\pi} \sqrt{\frac{z e^t}{(4 - z e^t)^{3/2}}} \left(\sin^{-1} \frac{\sqrt{z} e^{t/2}}{2} \cos^{-1} \frac{\sqrt{z}}{2} - \cos^{-1} \frac{\sqrt{z} e^{t/2}}{2} \sin^{-1} \frac{\sqrt{z}}{2} \right), \quad (16)$$

$$\rho_2(z, t) = 4\sqrt{\frac{(4-z)e^t}{(4-ze^t)^3}} \quad (17)$$

are easily verified. In the case $z = 2$, (15) and (16) are identical to Dyson's formulas [3, 5] obtained empirically.

In view of the prominent role that the ratio $R_1(z, t)/R_2(z, t)$ plays in Dyson et al. [3] for $z = 2$ it is interesting to examine it for general z . From (17) we have

$$R_2(z, k) = \frac{2k!\sqrt{4-z}}{\pi i} \oint \frac{ds}{s^{k+1}} \frac{e^{s/2}}{(4-ze^s)^{3/2}}. \quad (18)$$

The non-zero singularity closest to $s = 0$ is $s_0 = \ln(4/z)$ and it dominates the asymptotic behavior. Ignoring the other singularities, distorting the contour to a small circle about s_0 and translating back to the origin by $t = s - s_0$, we have

$$R_2(z, k) \sim -\frac{k!\sqrt{4-z}}{zs_0^{k+1}} \oint \frac{dt}{2\pi i} \frac{e^{t/2}}{(1-e^t)^{3/2}}. \quad (19)$$

The exact value of the integral in (19) is $-(2/\pi)\sqrt{e/(e-1)}$, and so

$$R_2(z, k) \sim \frac{k!}{(\ln(4/z))^{k+1}} \frac{2}{\pi} \sqrt{\frac{e(4-z)}{z(e-1)}}. \quad (20)$$

In the same way we obtain

$$R_1(z, k) \sim \frac{k!}{(\ln(4/z))^{k+1}} \left(\sqrt{2} + \frac{2}{\pi} \left(\sqrt{\frac{e}{e-1}} - \sqrt{2} \right) \cos^{-1} \frac{\sqrt{z}}{2} - \frac{2^{3/2}}{\pi} \sin^{-1} \frac{\sqrt{z}}{2} \right). \quad (21)$$

4 Discussion

From (20) and (21) we find

$$\lim_{k \rightarrow \infty} \left(\frac{R_1(z, k)}{R_2(z, k)} - \sqrt{\frac{z}{4-z}} \sin^{-1} \frac{\sqrt{z}}{2} \right) = \sqrt{\frac{z}{4-z}} \left(\cos^{-1} \frac{\sqrt{z}}{2} - \sin^{-1} \frac{\sqrt{z}}{2} \right). \quad (22)$$

It thus appears that Lehmer's choice, $z = 2$, is the unique permissible case for which the limit (22) vanishes. (Also the *Lehmer limit*, as defined by Dyson et al. [3], relates to $\pi/4$ here rather than π). Finally, for negative integer indices, since

$$2S_{-k}(z) = {}_{k+1}F_k(1, \dots, 1; \frac{3}{2}, 2, \dots, 2; \frac{1}{4}z), \quad (23)$$

the fact that $S_{-k}(z)$ can be obtained from $S_0(z)$ by successive integrations with respect to z and the explicit evaluations by Lehmer [2], Borwein and Girgensohn [4] and others [6, 7, 8, 9, 10] it should be possible to obtain explicit values for sundry generalized hypergeometric functions.

5 Appendix: Derivation of Equations (8) and (10)

Let us consider, for any integrable function F ,

$$I = \int_0^1 \frac{dt}{t} F(t(1-t))$$

Let $u = t(1-t)$, so $u(0) = u(1) = 0$; $u(1/2) = 1/4$. Then there are two expressions for t :

$$t_+ = \frac{1}{2}(1 + \sqrt{1-4t}) \quad \text{for } \frac{1}{2} \leq t \leq 1, \quad \text{with } \frac{dt_+}{t_+} = \left(1 - \frac{1}{\sqrt{1-4u}}\right) \frac{du}{u}$$

and

$$t_- = \frac{1}{2}(1 - \sqrt{1-4t}) \quad \text{for } 0 \leq t \leq \frac{1}{2}, \quad \text{with } \frac{dt_-}{t_-} = \left(1 + \frac{1}{\sqrt{1-4u}}\right) \frac{du}{u}.$$

Consequently,

$$\begin{aligned} I &= \int_0^{1/2} \frac{dt_-}{t_-} F(u) + \int_{1/2}^1 \frac{dt_+}{t_+} F(u) = 2 \int_0^{1/4} \frac{du}{u\sqrt{1-4u}} F(u) \\ &= 2 \int_0^1 \frac{dx}{x\sqrt{1-x}} F\left(\frac{1-x}{4}\right) = 2 \int_0^1 \frac{dt}{(1-t)\sqrt{t}} F\left(\frac{1-t}{4}\right) \end{aligned}$$

and, with $t = x^2$,

$$I = 4 \int_0^1 \frac{dx}{1-x^2} F\left(\frac{1-x^2}{4}\right).$$

Therefore,

$$L = \int_0^1 \frac{dt}{t} \frac{(zt(1-t))^\alpha}{(1-zt(1-t))^\beta} = 2 \left(\frac{z}{4}\right)^{\alpha-\beta} \int_0^1 dx \frac{(1-x^2)^{\alpha-1}}{(a^2+x^2)^\beta},$$

where $a^2 = 1/X = (4-z)/z$.

From standard references

$$\int_0^1 dx \cos(xy) (1-x^2)^{\alpha-1} = \sqrt{\frac{\pi y}{8}} \left(\frac{2}{y}\right)^\alpha \Gamma(\alpha) J_{\alpha-1/2}(y),$$

$$\int_0^\infty dx \cos(xy) (a^2+x^2)^{-\beta} = \frac{\sqrt{\pi}}{\Gamma(\beta)} \left(\frac{y}{2a}\right)^{\beta-1/2} K_{\beta-1/2}(ay)$$

so, by the Parseval relation for the Fourier transform

$$\int_0^1 dx \frac{(1-x^2)^{\alpha-1}}{(a^2+x^2)^\beta} = \frac{2^{\alpha-\beta} \Gamma(\alpha)}{a^{\beta-1/2} \Gamma(\beta)} \int_0^\infty dy y^{\beta-\alpha} J_{\alpha-1/2}(y) K_{\beta-1/2}(ay).$$

This is a tabulated Hankel Transform and yields

$$L = \sqrt{\pi} \left(\frac{4}{z}\right)^{\beta-\alpha} X^\beta \frac{\Gamma(\alpha)}{\Gamma(\alpha+1)} {}_2F_1\left(\frac{1}{2}, \beta; \alpha + \frac{1}{2}; -X\right).$$

Consequently

$$\int_0^1 \frac{dx}{x} \frac{(zt(1-t))^n}{(1-zt(1-t))^{n+1}} = \sqrt{\pi} \left(\frac{4}{z}\right) X^{n+1} {}_2F_1\left(\frac{1}{2}, n+1; n+\frac{1}{2}; -X\right)$$

However, since ${}_2F_1(a, b; c; z) = (1-z)^{c-a-b} {}_2F_1(c-a, c-b; c; z)$,

$${}_2F_1\left(\frac{1}{2}, n+1; n+\frac{1}{2}; -X\right) = (1+X)^{-1} {}_2F_1\left(-\frac{1}{2}; n; n+\frac{1}{2}; -X\right)$$

Next, we note that [11, p . 590]

$${}_2F_1(-1/2, 1; 3/2; z) = \frac{1}{2} \left(1 + (1-z) \frac{\tanh^{-1} \sqrt{z}}{\sqrt{z}}\right).$$

With $z \rightarrow -z$, noting that $-i \tanh^{-1} iw = \sin^{-1} \sqrt{\frac{w}{1+w}}$ one has

$${}_2F_1(-1/2, 1; 3/2; -z) = \frac{1}{2} (1 + (1+z) \frac{\sin^{-1} \sqrt{\frac{z}{1+z}}}{\sqrt{z}}). \quad (24)$$

We next apply Gauss' differentiation formula

$$\frac{d}{dz} ((1+z)^k {}_2F_1(-1/2, k; k+1/2; -z)) =$$

$$\frac{2k(k+1)}{2k+1} (1+z) {}_2F_1(-1/2, k+1; k+3/2; -z). \quad (25)$$

Iteration of (25) starting with (24), after a great deal of tedious algebra, aided by Mathematica, results in (10).

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