

HILBERT TRANSFORMS AND SUM RULES OF BESSEL MOMENTS

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To the memory of Jonathan M. Borwein (1951–2016)

ABSTRACT. Using Hilbert transforms, we establish two families of sum rules involving Bessel moments, which are integrals associated with Feynman diagrams in two-dimensional quantum field theory. With these linear relations among Bessel moments, we verify and generalize two conjectures by Bailey–Borwein–Broadhurst–Glasser and Broadhurst–Mellit.

1. INTRODUCTION

Let I_0 and K_0 be the modified Bessel functions of zeroth order, defined through the following Schläfli integral representations [19, §6.15]:

$$I_0(t) = \frac{1}{\pi} \int_0^\pi e^{t \cos \theta} d\theta, \quad K_0(t) = \int_0^\infty e^{-t \cosh u} du. \quad (1.1)$$

Bessel moments, in the form of absolutely convergent integrals

$$\mathbf{IKM}(a, b; c) := \int_0^\infty [I_0(t)]^a [K_0(t)]^b t^c dt \quad (1.2)$$

for certain non-negative integers a, b, c , arise naturally from perturbative diagrammatic expansions in two-dimensional quantum field theory [10, 1, 5, 9].

The Bessel moments $\mathbf{IKM}(a, b; c)$ involving no more than four Bessel functions (that is, $a + b \leq 4$) have been thoroughly studied [1]: arithmetically speaking, the $\mathbf{IKM}(a, b; 1)$ values are expressible as rational multiples of special L -values when $a + b \leq 4$; combinatorially speaking, generating functions for the corresponding sequences $\mathbf{IKM}(a, b; c)$ with respect to $c \in \mathbb{Z}_{>0}$ are explicitly known. In contrast, for $a + b \geq 5$, the combinatorial structure of Bessel moments $\mathbf{IKM}(a, b; c)$ and their relation to L -functions largely remain elusive. For the $a + b = 5$ case alone, closed-form evaluations of $\mathbf{IKM}(2, 3; 1)$ and $\mathbf{IKM}(1, 4; 1)$ have drawn heavily on critical advances in symbolic computation [1], algebraic geometry [3] and number theory [14, 16] within the last decade. A recent progress towards the understanding of Bessel moments with $a + b = 8, c = 1$ (see [5, §7.6] and [9, §7]) has benefited from Yun’s insights [20] into the Langlands program.

On a different note, numerical experimentations seem to support the existence of various algebraic relations among Bessel moments with $a + b \geq 5$. As examples of sum rules suggested from high-precision numerical computations, we mention here two open problems concerning Bessel moments: a cancelation formula proposed by Bailey–Borwein–Broadhurst–Glasser in 2008, and a conjecture on integrality formulated by Broadhurst (jointly with Mellit, and in honor of Crandall) in 2016.

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Conjecture 1.1 (B³G sum rule [1, “final conjecture”, (220)]). *For each pair of integers $(n, k) \in \mathbb{Z}^2$ satisfying $n \geq 2k \geq 2$, the following sum of integrals vanishes identically:*

$$Z_{2n, n-2k} := \sum_{m=0}^{\lfloor n/2 \rfloor} (-1)^m \binom{n}{2m} \int_0^\infty [\pi I_0(t)]^{n-2m} [K_0(t)]^{n+2m} t^{n-2k} dt = 0, \quad (1.3)$$

where $\binom{n}{j} = \frac{n!}{j!(n-j)!}$ and $\lfloor x \rfloor$ represents the greatest integer less than or equal to x .

Conjecture 1.2 (Broadhurst–Mellit integer sequence [5, (149) in Conjecture 5], *viz.* Crandall numbers). *For each $n \in \mathbb{Z}_{>0}$, the following integral*

$$A(n) := \left(\frac{2}{\pi}\right)^4 \int_0^\infty \{[\pi I_0(t)]^2 - [K_0(t)]^2\} I_0(t) [K_0(t)]^5 (2t)^{2n-1} dt \quad (1.4)$$

evaluates to an integer.

In this note, we harness the Hilbert transform to verify both Conjectures 1.1 and 1.2, in a unified and coherent framework, without explicitly evaluating individual Bessel moments contained in the sum rules.

The article is organized as follows. In §2, we start from a brief overview of the classical Hilbert transform (§2.1), a major analytic and algebraic tool in this work; we then compute Hilbert transforms for some simple algebraic expressions involving modified Bessel functions (§2.2), paving the way for the subsequent proofs of Conjecture 1.1 in §3.1, and Conjecture 1.2 in §3.2. In addition to proving both these conjectures, we construct a conjugate to the B³G sum rule in §3.1:

$$Y_{2n, n-2k} := \sum_{m=1}^{\lfloor n/2 \rfloor + 1} (-1)^m \binom{n}{2m-1} \int_0^\infty [\pi I_0(t)]^{n-2m+1} [K_0(t)]^{n+2m-1} t^{n-2k-1} dt = 0, \quad (1.5)$$

for $n-1 \geq 2k \geq 2$, and obtain generalizations of the Broadhurst–Mellit integer sequence in §3.2:

$$\frac{4}{\pi^{2m+1}} \sum_{\ell=1}^{m+1} (-1)^{\ell-1} \binom{2m}{2\ell-1} \int_0^\infty [\pi I_0(t)]^{2(m-\ell)+1} [K_0(t)]^{2(m+\ell)-1} (2t)^{2(n+m)-3} dt \in \mathbb{Z}_{>0}, \quad (1.6)$$

$$\frac{4^n}{\pi^{2m}} \sum_{\ell=1}^m (-1)^{\ell-1} \binom{2m-1}{2\ell-1} \int_0^\infty [\pi I_0(t)]^{2(m-\ell)} [K_0(t)]^{2(m+\ell-1)} (2t)^{2(n+m-2)} dt \in \mathbb{Z}_{>0}, \quad (1.7)$$

for all $m, n \in \mathbb{Z}_{>0}$.

2. HILBERT TRANSFORMS INVOLVING MODIFIED BESSEL FUNCTIONS

2.1. Basic properties of Hilbert transforms. The key device in our proof is the Hilbert transform $\widehat{\mathcal{H}}$, which operates on a suitably regular function $f(x)$, a.e. $x \in \mathbb{R}$ through a Cauchy principal value:¹

$$(\widehat{\mathcal{H}}f)(x) = \mathcal{P} \int_{-\infty}^\infty \frac{f(\xi) d\xi}{\pi(x-\xi)}, \quad \text{a.e. } x \in \mathbb{R}. \quad (2.1)$$

When restricted to an L^p space for $1 < p < \infty$, the Hilbert transform induces a bounded linear operator $\widehat{\mathcal{H}} : L^p(\mathbb{R}) \rightarrow L^p(\mathbb{R})$ [18, p. 188].

We list here three fundamental properties of the Hilbert transform on \mathbb{R} :

¹For our purposes, it suffices to evaluate the Hilbert transform of a function for almost every (a.e.) real variable, leaving out a subset of measure zero that will not affect subsequent computations of Lebesgue integrals. Later afterwards, an equal sign may also be used to denote an equality that is valid almost everywhere (even when “a.e.” is not written), depending on context.

(HT1) (**Parseval-type Identity**) For $f \in L^p(\mathbb{R}), p > 1; g \in L^q(\mathbb{R}), q > 1$ and $\frac{1}{p} + \frac{1}{q} = 1$, we have [11, (10.81)]:

$$\langle f, \widehat{\mathcal{H}}g \rangle + \langle g, \widehat{\mathcal{H}}f \rangle = 0, \quad \text{where } \langle f_1, f_2 \rangle := \int_{-\infty}^{\infty} f_1(x)f_2(x)dx. \quad (2.2)$$

(HT2) (**Hardy–Poincaré–Bertrand Formula**) For $f \in L^p(\mathbb{R}), p > 1; g \in L^q(\mathbb{R}), q > 1$ and $\frac{1}{p} + \frac{1}{q} = 1$, we have [11, (2.227)]:

$$\widehat{\mathcal{H}}(f \widehat{\mathcal{H}}g + g \widehat{\mathcal{H}}f) = (\widehat{\mathcal{H}}f)(\widehat{\mathcal{H}}g) - fg. \quad (2.3)$$

(HT3) (**Moment Formula**) For $n \in \mathbb{Z}_{>0}$, the following identity [11, (4.113)]

$$\mathcal{P} \int_{-\infty}^{\infty} \frac{\xi^n f(\xi) d\xi}{\pi(x-\xi)} = x^n \mathcal{P} \int_{-\infty}^{\infty} \frac{f(\xi) d\xi}{\pi(x-\xi)} - \frac{1}{\pi} \sum_{k=0}^{n-1} x^k \int_{-\infty}^{\infty} \xi^{n-1-k} f(\xi) d\xi \quad (2.4)$$

holds when each term therein is well-defined.

Due to frequent invocations of properties (HT1) and (HT2), we will abbreviate the relations in (2.2) and (2.3) as $\text{Par}(f, g)$ and $\text{HPB}(f, g)$, respectively.

As a consequence of the Hardy–Poincaré–Bertrand formula, the operator $i \widehat{\mathcal{H}}$ is its own inverse [11, (4.18)]. In other words,

$$\widehat{\mathcal{H}}^2 = -\widehat{I}: L^p(\mathbb{R}) \longrightarrow L^p(\mathbb{R}), \quad 1 < p < \infty \quad (2.5)$$

maps an L^p function $f(x), x \in \mathbb{R}$ to $-f(x), x \in \mathbb{R}$.

2.2. Applications to modified Bessel functions. To simplify notations, we introduce short-hands for certain expressions involving the modified Bessel functions.

Definition 2.1. For a real variable $x \in \mathbb{R}$, we define the functions ι, ι_+, ι_- , and $\kappa, \kappa_+, \kappa_-$ as follows:

$$\iota(x) := \pi I_0(x); \quad \kappa(x) := K_0(|x|); \quad (2.6)$$

$$\iota_+(x) := \iota(x)e^{-x} \mathbb{1}_{(0, \infty)}(x); \quad \kappa_+(x) := \kappa(x)e^{-x}; \quad (2.7)$$

$$\iota_-(x) := \iota(x)e^x \mathbb{1}_{(-\infty, 0)}(x); \quad \kappa_-(x) := \kappa(x)e^x, \quad (2.8)$$

where the indicator function behaves like $\mathbb{1}_A(x) = 1, x \in A; \mathbb{1}_A(x) = 0, x \notin A$. It is clear that $\kappa, \iota \in L^p(\mathbb{R})$ for $p > 1$ and $\iota_+, \iota_-, \kappa_+, \kappa_- \in L^p(\mathbb{R})$ for $p > 2$.

Lemma 2.2 (Hilbert transform of modified Bessel functions). *We have the following Hilbert transform formulae:*

$$\widehat{\mathcal{H}} \iota_+ = -\kappa_+, \quad \widehat{\mathcal{H}} \iota_- = \kappa_-, \quad (2.9)$$

$$\widehat{\mathcal{H}} \kappa_+ = \iota_+, \quad \widehat{\mathcal{H}} \kappa_- = -\iota_-, \quad (2.10)$$

which entail the following Feynman rules:

$$\widehat{\mathcal{H}}(\iota \kappa \text{sgn}) = -\kappa^2, \quad (2.11)$$

$$\widehat{\mathcal{H}}(\kappa^2) = \iota \kappa \text{sgn}, \quad (2.12)$$

for $\text{sgn}(x) = x/|x|, \forall x \in \mathbb{R} \setminus \{0\}$.

Proof. The results in (2.9) and (2.10) are classical (see [11, §9.10] or [12, Appendix 1, Tables 1.8H and 1.8I]). Furthermore, we note that (2.10) follows from (2.9) and (2.5).

Specializing the Hardy–Poincaré–Bertrand formula to $\text{HPB}(\iota_+, \iota_- \mathbb{1}_{[-L, L]})$ where the values of $\iota_- \mathbb{1}_{[-L, L]}$ are zero-padded outside a bounded interval $[-L, L]$, and exploiting the continuity of the Hilbert transform as a linear operator, we obtain the following formula in the $L \rightarrow \infty$ limit:

$$\widehat{\mathcal{H}}(\iota_+ \kappa_- - \iota_- \kappa_+) = -\kappa_+ \kappa_- - \iota_+ \iota_-. \quad (2.13)$$

Here, the left-hand side of the equation above is equal to $\widehat{\mathcal{H}}(\iota\kappa \operatorname{sgn})$, while the right-hand side simplifies to $-\kappa^2$. This proves (2.11). Applying $\widehat{\mathcal{H}}$ to both sides of (2.11), we arrive at (2.12). \blacksquare

Remark With a Parseval-type identity $\operatorname{Par}(\iota\kappa \operatorname{sgn}, \kappa^2)$, we immediately deduce from (2.11) and (2.12) a sum rule for Bessel moments [1, (91) and (219)]

$$\begin{aligned} 0 &= \frac{\langle \iota\kappa \operatorname{sgn}, \widehat{\mathcal{H}}(\kappa^2) \rangle + \langle \kappa^2, \widehat{\mathcal{H}}(\iota\kappa \operatorname{sgn}) \rangle}{2} \\ &= \int_0^\infty [\pi I_0(t)]^2 [K_0(t)]^2 dt - \int_0^\infty [K_0(t)]^4 dt =: Z_{4,0}. \end{aligned} \quad (2.14)$$

Unlike the analysis in [1, (91)], during our proof of the cancelation formula above, we have not explicitly computed either $\int_0^\infty [\pi I_0(t)]^2 [K_0(t)]^2 dt$ or $\int_0^\infty [K_0(t)]^4 dt$. \square

3. APPLICATIONS TO SUM RULES OF BESSEL MOMENTS

3.1. B³G sum rule. To the best of our knowledge, except for the case of $Z_{4,0}$ (which was solved by brute force [1, (91)] and reaffirmed in §2.2 above), Conjecture 1.1 has remained hitherto open. Both [1, (221)] and [5, (112)] have singled out the verification of $Z_{6,1} = 0$ as an unsolved problem.²

Before handling Conjecture 1.1 in full, we first walk through the proof of $Z_{6,1} = 0$ and $Z_{8,0} = 0$ in the next lemma, to illustrate our strategies. In what follows, for $n \in \mathbb{Z}_{>0}$, we set $\varpi_n(x) := x^n, \forall x \in \mathbb{R}$; we further define $\varpi_0(x) := 1, \forall x \in \mathbb{R}$.

Lemma 3.1. *We have $Z_{6,1} = 0$ and $Z_{8,0} = 0$.*

Proof. By $\operatorname{HPB}(\iota\kappa \operatorname{sgn}, \kappa^2)$, we have a Hilbert transform formula:

$$\widehat{\mathcal{H}}(\iota^2 \kappa^2 - \kappa^4) = -2\iota\kappa^3 \operatorname{sgn}. \quad (3.1)$$

Appealing to (2.4), we deduce

$$\widehat{\mathcal{H}}((\iota^2 \kappa^2 - \kappa^4)\varpi_1) = -2\iota\kappa^3 \varpi_1 \operatorname{sgn} - \frac{1}{\pi} \int_{-\infty}^\infty (\iota^2 \kappa^2 - \kappa^4)(\xi) d\xi = -2\iota\kappa^3 \varpi_1 \operatorname{sgn}, \quad (3.2)$$

upon invoking $Z_{4,0} = 0$ in the last step. By $\operatorname{Par}(\kappa^2, (\iota^2 \kappa^2 - \kappa^4)\varpi_1)$, we arrive at

$$\begin{aligned} 0 &= \frac{\langle \kappa^2, -2\iota\kappa^3 \varpi_1 \operatorname{sgn} \rangle + \langle \iota\kappa \operatorname{sgn}, (\iota^2 \kappa^2 - \kappa^4)\varpi_1 \rangle}{2} \\ &= \int_0^\infty [\pi I_0(t)]^3 [K_0(t)]^3 t dt - 3 \int_0^\infty \pi I_0(t) [K_0(t)]^5 t dt = Z_{6,1}, \end{aligned} \quad (3.3)$$

as claimed.

By (3.1) and $\operatorname{Par}(\iota^2 \kappa^2 - \kappa^4, 2\iota\kappa^3 \operatorname{sgn})$, we have

$$\begin{aligned} 0 &= \frac{\langle \iota^2 \kappa^2 - \kappa^4, \widehat{\mathcal{H}}(2\iota\kappa^3 \operatorname{sgn}) \rangle + \langle 2\iota\kappa^3 \operatorname{sgn}, \widehat{\mathcal{H}}(\iota^2 \kappa^2 - \kappa^4) \rangle}{2} \\ &= \int_0^\infty [\pi I_0(t)]^4 [K_0(t)]^4 dt - 6 \int_0^\infty [\pi I_0(t)]^2 [K_0(t)]^6 dt + \int_0^\infty [K_0(t)]^8 dt = Z_{8,0}, \end{aligned} \quad (3.4)$$

as stated. \blacksquare

In the arguments above, we have established the sum rules $Z_{6,1} = 0$ and $Z_{8,0} = 0$ on a simpler vanishing identity $Z_{4,0} = 0$. In the next proposition, we proceed to more general transitions from 2ℓ Bessel factors to $2\ell + 2$ and $2\ell + 4$ scenarios, where $\ell \in \mathbb{Z}_{>0}$.

²Confusingly, in footnote 13 of [1], the authors apparently declared that the conjecture on $Z_{2n, n-2k} = 0$ “has also now been resolved”, without supplying further citations to publications or preprints. This contradicts the recent claims in [5, (112)] and [9, (5.6)] on the open status of $Z_{6,1} = 0$.

Proposition 3.2 (Hilbert ladders). *We have the following Feynman rules for all $\ell \in \mathbb{Z}_{>0}$:*

$$\begin{cases} \widehat{\mathcal{H}}(\zeta_\ell) = \eta_\ell, \\ \widehat{\mathcal{H}}(\eta_\ell) = -\zeta_\ell, \end{cases} \quad (3.5)$$

where the Hilbert ladders are defined by

$$\begin{cases} 2\zeta_\ell = \kappa^\ell [(\iota \operatorname{sgn} + i\kappa)^\ell + (\iota \operatorname{sgn} - i\kappa)^\ell], \\ 2\eta_\ell = i\kappa^\ell [(\iota \operatorname{sgn} + i\kappa)^\ell - (\iota \operatorname{sgn} - i\kappa)^\ell]. \end{cases} \quad (3.6)$$

Moreover, we have

$$\widehat{\mathcal{H}}(\zeta_\ell \omega_j) = \eta_\ell \omega_j, \quad \forall j \in \mathbb{Z} \cap [0, \ell). \quad (3.7)$$

Proof. By direct computation, we can verify the following algebraic relations for Hilbert ladders:

$$\zeta_\ell \eta_m + \eta_\ell \zeta_m = \eta_{\ell+m}, \quad (3.8)$$

$$\zeta_\ell \zeta_m - \eta_\ell \eta_m = \zeta_{\ell+m}. \quad (3.9)$$

For $\ell = 1$, the Feynman rules in (3.5) reproduce (2.11)–(2.12) in Lemma 2.2. Assuming that the Hilbert transform formula $\widehat{\mathcal{H}}(\zeta_\ell) = \eta_\ell$ holds up to a certain integer ℓ , then $\operatorname{HPB}(\zeta_\ell, \zeta_1)$ and (3.8)–(3.9) bring us $\widehat{\mathcal{H}} \eta_{\ell+1} = -\zeta_{\ell+1}$, which is equivalent to $\widehat{\mathcal{H}}(\zeta_{\ell+1}) = \eta_{\ell+1}$. By induction, this proves (3.5) in its entirety.

For $\ell = 1$, the statement in (3.7) is just $\widehat{\mathcal{H}}(\zeta_1) = \eta_1$. For $\ell = 2$, the relation $\widehat{\mathcal{H}}(\zeta_2 \omega_1) = \eta_2 \omega_1$ is effectively proved in (3.2). Now suppose that up to a certain integer ℓ , we have $\widehat{\mathcal{H}}(\zeta_\ell \omega_j) = \eta_\ell \omega_j, \forall j \in \mathbb{Z} \cap [0, \ell)$. From $\operatorname{Par}(\zeta_\ell \omega_j, \eta_2 \omega_1)$ and $\zeta_{\ell+2} \omega_{j+1} = (\zeta_\ell \omega_j)(\zeta_2 \omega_1) - (\eta_\ell \omega_j)(\eta_2 \omega_1) = -(\zeta_\ell \omega_j) \widehat{\mathcal{H}}(\eta_2 \omega_1) - (\eta_2 \omega_1) \widehat{\mathcal{H}}(\zeta_\ell \omega_j)$, we know that $\int_{\mathbb{R}} \zeta_{\ell+2}(x) x^{j+1} dx = 0, \forall j \in \mathbb{Z} \cap [0, \ell)$. According to the moment formula in (2.4), it follows that $\widehat{\mathcal{H}}(\zeta_{\ell+2} \omega_{j+2}) = \eta_{\ell+2} \omega_{j+2}, \forall j \in \mathbb{Z} \cap [0, \ell)$. Additionally, $\operatorname{Par}(\zeta_\ell, \eta_2)$ and $\zeta_{\ell+2} = \zeta_\ell \zeta_2 - \eta_\ell \eta_2 = -\zeta_\ell \widehat{\mathcal{H}} \eta_2 - \eta_2 \widehat{\mathcal{H}} \zeta_\ell$ tell us that $\int_{\mathbb{R}} \zeta_{\ell+2}(x) dx = 0$. Thus $\widehat{\mathcal{H}}(\zeta_{\ell+2} \omega_1) = \eta_{\ell+2} \omega_1$ is also true. This proves that $\widehat{\mathcal{H}}(\zeta_{\ell+2} \omega_j) = \eta_{\ell+2} \omega_j, \forall j \in \mathbb{Z} \cap [0, \ell+2)$, so (3.7) follows by induction. ■

Theorem 3.3 (B³G sum rule). *We have*

$$Z_{2n, n-2k} := \sum_{m=0}^{\lfloor n/2 \rfloor} (-1)^m \binom{n}{2m} \int_0^\infty [\pi I_0(t)]^{n-2m} [K_0(t)]^{n+2m} t^{n-2k} dt = 0 \quad (3.10)$$

for all integer pairs (n, k) meeting the requirements $n \geq 2k \geq 2$.

Proof. Simply notice that binomial expansion leads us to

$$\zeta_n = \sum_{m=0}^{\lfloor n/2 \rfloor} (-1)^m \binom{n}{2m} (\iota \operatorname{sgn})^{n-2m} \kappa^{n+2m}, \quad (3.11)$$

and (3.7) implies that $\int_{\mathbb{R}} \zeta_n(x) x^j dx = 0$ for $j \in \mathbb{Z} \cap [0, n-1)$. ■

Remark Since Hilbert inversion brings us $\widehat{\mathcal{H}}(\eta_\ell \omega_j) = -\zeta_\ell \omega_j = \omega_j \widehat{\mathcal{H}} \eta_\ell$ for $j \in \mathbb{Z} \cap [0, \ell)$, we must have $\int_{\mathbb{R}} \eta_\ell(x) x^j dx = 0$ for $j \in \mathbb{Z} \cap [0, \ell-1)$, as well. Excluding from these vanishing identities the trivial cases where the integrands are odd functions over \mathbb{R} , we arrive at another family of cancelation formulae, conjugate to the B³G sum rule:

$$Y_{2n, n-2k} := \sum_{m=1}^{\lfloor n/2 \rfloor + 1} (-1)^m \binom{n}{2m-1} \int_0^\infty [\pi I_0(t)]^{n-2m+1} [K_0(t)]^{n+2m-1} t^{n-2k-1} dt = 0, \quad (3.12)$$

where $n-1 \geq 2k \geq 2$. □

Remark Before writing this paper, we constructed alternative (and actually simpler) proofs of the B^3G sum rule (3.10) and its conjugate (3.12), via a vanishing contour integral

$$\int_{-i\infty}^{i\infty} z^m [H_0^{(1)}(z)H_0^{(2)}(z)]^n dz = 0, \quad m \in \mathbb{Z} \cap [0, n-1], \quad (3.13)$$

where $H_0^{(1)}, H_0^{(2)}$ are cylindrical Hankel functions and the contour can be closed to the right [cf. 19, §7.2]. In spite of this complex analytic shortcut, we encountered “infrared divergence” (singular behavior of the integrand as $|z| \rightarrow \infty$) in (3.13), when we attempted to raise the power m to higher values and calculate $\mathbf{IKM}(a, b; m)$ for $m+1 \geq (a+b)/2$, which is the situation occurring in the Broadhurst–Mellit integer sequence (see §3.2 below). As we will soon see, the Hilbert transform method will succeed where the contour integration approach fails. \square

3.2. Broadhurst–Mellit integer sequence (Crandall numbers). Let

$$A(n) := (2/\pi)^4 2^{2n-1} [\pi^2 \mathbf{IKM}(3, 5; 2n-1) - \mathbf{IKM}(1, 7; 2n-1)], \quad \forall n \in \mathbb{Z}_{>0} \quad (3.14)$$

be the sequence defined in (1.4). Despite (erstwhile) conjectural integrality, this now becomes entry A262961 in the On-line Encyclopedia of Integer Sequences (OEIS) [17, under the title “Crandall numbers”, in homage to Richard E. Crandall (1947–2012), who made seminal contributions to the computation of Bessel moments], along with the following comments (at the time of writing, in April 2017):

*“Anton Mellit and David Broadhurst define the sequence to be the ‘**round**’ (emphasis added) of the integral, with the conjecture that this rounding is exact. No one seems to know how to prove that any of the integrals gives a rational number, let alone an integer.”*

By (3.12), we see that $Y_{8,2} = 0$ implies the following vanishing identity

$$A(1) := \frac{32}{\pi^4} \int_0^\infty \{[\pi I_0(t)]^2 - [K_0(t)]^2\} I_0(t)[K_0(t)]^5 t dt = 0, \quad (3.15)$$

which was first conjectured by Broadhurst in [5, (148)], after collaboration with Mellit. The rest of the sequence $A(n+1), n \in \mathbb{Z}_{>0}$ in Conjecture 1.2 is characterized by the theorem below.

Theorem 3.4 (Broadhurst–Mellit integer sequence, *viz.* Crandall numbers). *For all $n \in \mathbb{Z}_{>1}$, we have*

$$A(n) := \left(\frac{2}{\pi}\right)^4 \int_0^\infty \{[\pi I_0(t)]^2 - [K_0(t)]^2\} I_0(t)[K_0(t)]^5 (2t)^{2n-1} dt \in \mathbb{Z}_{>0}. \quad (3.16)$$

Furthermore, the following explicit formula holds for $n \in \mathbb{Z}_{>0}$:

$$A(n+1) = \frac{1}{2^{4(n-1)}} \sum_{m=1}^n \sum_{\ell=1}^m \sum_{k=1}^{\ell} \frac{[(2n-2m)!]^3 [(2m-2\ell)!]^3 [(2\ell-2k)!]^3 [(2k-2)!]^3}{[(n-m)!]^4 [(m-\ell)!]^4 [(\ell-k)!]^4 [(k-1)!]^4}. \quad (3.17)$$

Proof. We recall the following Bessel moments from [1, (7)] (see also [13, 4]):

$$\int_0^\infty [K_0(t)]^2 t^{2n} dt = \frac{\pi^2 [(2n)!]^3}{4^{3n+1} (n!)^4}, \quad \forall n \in \mathbb{Z}_{\geq 0}, \quad (3.18)$$

which enable us to compute the following Hilbert transforms with the aid of (2.4) and (2.12):

$$\begin{aligned} \widehat{\mathcal{H}}(\kappa^2 \omega_{2n}) &= \omega_{2n} \widehat{\mathcal{H}}(\kappa^2) - \frac{1}{\pi} \sum_{k=0}^{2n-1} \omega_k \int_{-\infty}^{\infty} \xi^{2n-1-k} [\kappa(\xi)]^2 d\xi \\ &= i\kappa \omega_{2n} \operatorname{sgn} - \frac{\pi}{2} \sum_{m=1}^n \frac{[(2n-2m)!]^3}{4^{3(n-m)} [(n-m)!]^4} \omega_{2m-1}, \quad \forall n \in \mathbb{Z}_{>0}. \end{aligned} \quad (3.19)$$

Spelling out $\text{Par}(-3\iota^2\kappa^4\omega_1 + \kappa^6\omega_1, \kappa^2\omega_{2n})$ [i.e. $\text{Par}(\eta_3\omega_1, \kappa^2\omega_{2n})$], we see that $A(n+1), \forall n \in \mathbb{Z}_{>0}$ is expressible as a finite sum of Bessel moments involving only six Bessel functions:

$$A(n+1) = \frac{2^{2(n+1)}}{\pi^4} \sum_{m=1}^n \frac{[(2n-2m)!]^3}{4^{3(n-m)}[(n-m)!]^4} \int_0^\infty \{3[\pi I_0(t)]^2 - [K_0(t)]^2\} [K_0(t)]^4 t^{2m} dt. \quad (3.20)$$

By $\text{Par}(2\iota\kappa^3 \text{sgn}, \kappa^2\omega_{2m})$, we arrive at further reduction:

$$\begin{aligned} & \int_0^\infty \{3[\pi I_0(t)]^2 - [K_0(t)]^2\} [K_0(t)]^4 t^{2m} dt \\ &= \pi^2 \sum_{\ell=1}^m \frac{[(2m-2\ell)!]^3}{4^{3(m-\ell)}[(m-\ell)!]^4} \int_0^\infty I_0(t) [K_0(t)]^3 t^{2\ell-1} dt. \end{aligned} \quad (3.21)$$

With $\widehat{\mathcal{H}}(\kappa^2\omega_1) = \iota\kappa\omega_1 \text{sgn}$ and $\text{Par}(\kappa^2\omega_1, \kappa^2\omega_{2(\ell-1)})$, we evaluate the remaining Bessel moments involving four Bessel functions:

$$\int_0^\infty I_0(t) [K_0(t)]^3 t^{2\ell-1} dt = \frac{\pi^2}{4^{3\ell-1}} \sum_{k=1}^{\ell} \frac{[(2\ell-2k)!]^3 [(2k-2)!]^3}{[(\ell-k)!]^4 [(k-1)!]^4}. \quad (3.22)$$

So far, we have verified (3.17).

Next, we show that

$$\alpha_\ell := \frac{1}{2^{4(\ell-1)}} \sum_{k=1}^{\ell} \frac{[(2\ell-2k)!]^3 [(2k-2)!]^3}{[(\ell-k)!]^4 [(k-1)!]^4} \in \mathbb{Z}, \quad \forall \ell \in \mathbb{Z}_{>0}. \quad (3.23)$$

Towards this end, we first recall from [1, (55) and (56)] the following formula

$$\int_0^\infty I_0(t) [K_0(t)]^3 t^{2\ell-1} dt = \frac{\pi^2}{16} \left[\frac{(\ell-1)!}{4^{\ell-1}} \right]^2 \sum_{k=0}^{\ell-1} \binom{\ell-1}{k}^2 \binom{2(\ell-1-k)}{\ell-1-k} \binom{2k}{k}, \quad (3.24)$$

which combines with (3.22) into³

$$\alpha_\ell = \frac{[(\ell-1)!]^2}{4^{\ell-1}} \sum_{k=0}^{\ell-1} \binom{\ell-1}{k}^2 \binom{2(\ell-1-k)}{\ell-1-k} \binom{2k}{k} = \frac{[(\ell-1)!]^2}{4^{\ell-1}} D_{\ell-1}, \quad (3.25)$$

where $D_j, j \in \mathbb{Z}_{\geq 0}$ is called the j -th Domb number in combinatorics. In [15, Theorem 3.1], Rogers has shown that the following identity holds for $|u|$ sufficiently small:

$${}_3F_2 \left(\begin{matrix} \frac{1}{3}, \frac{1}{2}, \frac{2}{3} \\ 1, 1 \end{matrix} \middle| \frac{27u^2}{4(1-u)^3} \right) = (1-u) \sum_{n=0}^{\infty} \frac{D_n}{4^n} u^n. \quad (3.26)$$

Here, the generalized hypergeometric series ${}_pF_q$ is defined as

$${}_pF_q \left(\begin{matrix} a_1, \dots, a_p \\ b_1, \dots, b_q \end{matrix} \middle| x \right) := \sum_{n=0}^{\infty} \frac{(a_1)_n \cdots (a_p)_n}{(b_1)_n \cdots (b_q)_n} \frac{x^n}{n!}, \quad (3.27)$$

where the Pochhammer symbol represents the rising factorial

$$(a)_n := \begin{cases} 1, & n = 0 \\ a(a+1)\cdots(a+n-1), & n \in \mathbb{Z}_{>0} \end{cases}; \quad (a)_n := \frac{\Gamma(a+n)}{\Gamma(a)}, \quad a \notin \mathbb{Z}_{\leq 0},$$

for $\{b_1, \dots, b_q\} \cap \mathbb{Z}_{\leq 0} = \emptyset$. In Rogers' identity (3.26), the coefficient of u^n for $n \in \mathbb{Z}_{>0}$ is equal to

$$\frac{D_n}{4^n} - \frac{D_{n-1}}{4^{n-1}}, \quad (3.28)$$

³Naïvely, upon observing that $(n!)^2/2^n \in \mathbb{Z}, \forall n \in \mathbb{Z}_{\geq 4}$ and $D_n \in \mathbb{Z}, \forall n \in \mathbb{Z}_{\geq 0}$, we obtain $2^{\ell-1}\alpha_\ell \in \mathbb{Z}, \forall \ell \in \mathbb{Z}_{>0}$, at best. The divisibility statement $4^n \mid (n!)^2 D_n, \forall n \in \mathbb{Z}_{\geq 0}$ is thus deeper than these naïve observations. In our proof of the integrality $\alpha_\ell \in \mathbb{Z}, \forall \ell \in \mathbb{Z}_{>0}$, we need Rogers' work on modular forms [15], which in turn, was inspired by Bertin's studies of modular parametrizations for certain families of Calabi–Yau manifolds [2].

according to the right-hand side. Meanwhile, the Taylor expansion of the left-hand side reads

$${}_3F_2\left(\begin{matrix} \frac{1}{3}, \frac{1}{2}, \frac{2}{3} \\ 1, 1 \end{matrix} \middle| \frac{27u^2}{4(1-u)^3}\right) = 1 + 3 \sum_{n=1}^{\infty} [(2n-1)!!]^2 \binom{3n-1}{2n} \frac{1}{(n!2^n)^2} \frac{u^{2n}}{(1-u)^{3n}}, \quad (3.29)$$

where the double factorial is defined as $(2n-1)!! := (2n)/(n!2^n) \in \mathbb{Z}$ for $n \in \mathbb{Z}_{>0}$. As we compute the contribution from the aforementioned Taylor coefficients to

$$(n!)^2 \left(\frac{D_n}{4^n} - \frac{D_{n-1}}{4^{n-1}} \right) = \alpha_{n+1} - n^2 \alpha_n, \quad (3.30)$$

we are gathering finitely many summands, each of which is an integer multiple of $(k!!)^2 \in \mathbb{Z}$ for a certain odd positive integer k less than n . Therefore, we have $\alpha_1 = 1$ and $\alpha_{\ell+1} - \ell^2 \alpha_\ell \in \mathbb{Z}$ for all $\ell \in \mathbb{Z}_{>0}$.

Finally, by discrete convolution, we see that $A(n+1)$ is the coefficient of x^{n-1} in the polynomial

$$\left(\sum_{\ell=1}^{n+1} \alpha_\ell x^{\ell-1} \right)^2 = \left[\sum_{\ell=0}^n \frac{(\ell!)^2 D_\ell}{4^\ell} x^\ell \right]^2 \in \mathbb{Z}[x], \quad (3.31)$$

so $A(n+1) \in \mathbb{Z}$ must hold. ■

Remark More generally, we define the positive integer $\alpha_n^{[m]}$ as the coefficient of x^{n-1} in the polynomial $(\sum_{\ell=1}^{n+1} \alpha_\ell x^{\ell-1})^m \in \mathbb{Z}[x]$, for every $m \in \mathbb{Z}_{>0}$. By repeated applications of the recursions for Hilbert ladders [see (3.8) and (3.9)], we can show that

$$\alpha_n^{[m]} = \frac{4}{\pi^{2m+1}} \sum_{\ell=1}^{m+1} (-1)^{\ell-1} \binom{2m}{2\ell-1} \int_0^\infty [\pi I_0(t)]^{2(m-\ell)+1} [K_0(t)]^{2(m+\ell)-1} (2t)^{2(n+m)-3} dt. \quad (3.32)$$

Thus, the Broadhurst–Mellit integer sequence $A(1) = 0, A(n+1) = \alpha_n^{[2]}, n \in \mathbb{Z}_{>0}$ is just a special case within an infinite family of (linear combinations for) Bessel moments, preceded by

$$\alpha_n \equiv \alpha_n^{[1]} = \frac{8}{\pi^2} \int_0^\infty I_0(t) [K_0(t)]^3 (2t)^{2n-1} dt \in \mathbb{Z}_{>0} \quad (3.33)$$

and followed by

$$\alpha_n^{[3]} = \frac{8}{\pi^6} \int_0^\infty \{[\pi I_0(t)]^2 - 3[K_0(t)]^2\} \{3[\pi I_0(t)]^2 - [K_0(t)]^2\} I_0(t) [K_0(t)]^7 (2t)^{2n+3} dt \in \mathbb{Z}_{>0}, \quad (3.34)$$

for $n \in \mathbb{Z}_{>0}$. □

Remark For $\mathbf{IKM}(a, b; c)$ satisfying $c \equiv 0 \pmod{2}$ and $a + b \equiv 2 \pmod{4}$, we also have a family of integer sequences that runs parallel to (3.32). With the observation that

$$\frac{[(2n)!]^3}{2^{4n}(n!)^4} = \frac{1}{2^{2n}} [(2n-1)!!]^2 \binom{2n}{n} \in \frac{1}{2^{2n}} \mathbb{Z}, \quad \forall n \in \mathbb{Z}_{\geq 0}, \quad (3.35)$$

we define $\beta_n^{[m]}$ as the coefficient of x^{n-1} in the polynomial

$$\left[\sum_{\ell=0}^n \frac{(\ell!)^2 D_\ell}{4^\ell} x^\ell \right]^{m-1} \sum_{k=0}^n \frac{[(2k)!]^3}{2^{4k}(k!)^4} x^k \in \frac{1}{2^{2(n-1)}} \mathbb{Z}[x], \quad (3.36)$$

for every $m \in \mathbb{Z}_{>0}$. According to the discrete convolutions of Bessel moments that descend from the recursions for Hilbert ladders, we have

$$\beta_n^{[m]} = \frac{4}{\pi^{2m}} \sum_{\ell=1}^m (-1)^{\ell-1} \binom{2m-1}{2\ell-1} \int_0^\infty [\pi I_0(t)]^{2(m-\ell)} [K_0(t)]^{2(m+\ell-1)} (2t)^{2(n+m-2)} dt \in \frac{1}{2^{2(n-1)}} \mathbb{Z}_{>0} \quad (3.37)$$

for $m, n \in \mathbb{Z}_{>0}$. We can even combine the foregoing statements about $\alpha_n^{[m]}$ and $\beta_n^{[m]}$ into a more compact form, as follows:

$$\frac{2^{1+2(n-1)[1-(-1)^M]}}{\pi^{M+1}} \int_0^\infty \frac{[\pi I_0(t) + iK_0(t)]^M - [\pi I_0(t) - iK_0(t)]^M}{i} [K_0(t)]^M (2t)^{2n+M-3} dt \in \mathbb{Z}_{>0} \quad (3.38)$$

for all $M, n \in \mathbb{Z}_{>0}$. This confirms a recent empirical observation by Broadhurst and Roberts [6, Conjecture 2]. \square

To summarize, we have demonstrated that many sum rules of Bessel moments indeed issue from the three fundamental laws (HT1)–(HT3) of Hilbert transforms. We remind our readers that the methods developed in this note do not exhaust all possible linear relations for Feynman diagrams in two-dimensional quantum field theory. For example, the sum rule $9\pi^2 \mathbf{IKM}(4, 4; 1) - 14 \mathbf{IKM}(2, 6; 1) = 0$ [conjectured in 5, (147)] is not a simple consequence of Hilbert transforms, and its recent verification in [21, §5.2] involves heavy use of contour integrations over modular forms.

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